

Steady-state and time-dependent LPP modelling



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Outline

1. A steady state plasma (optically thin)

4d-4f + 4p-4d emission: statistical UTA (gf), ion distribution $f_z = f(T_e)$
 n_e and pulse length (YAG and CO₂)

2. A time-dependent 1-D plasma (optically thick)

ion distribution, $f_z = f(r, t)$, hydrodynamics [Medusa], some diagnostics
 N_i , N_j level populations for radiation transport [UCD],
2 surveys: power density and pulse length (optimum laser conditions)

3. A 2-D plasma

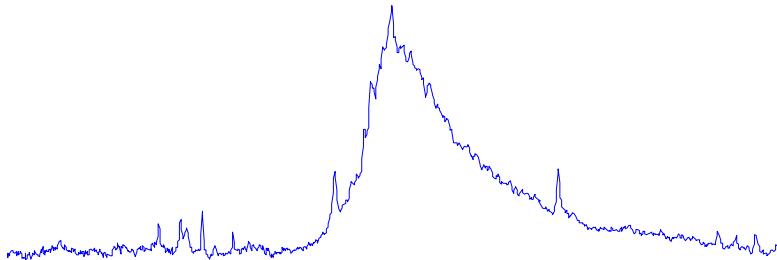
time-dependent spectra [EPPRA, Z*]
spatial pulse shape (flat-top/Gaussian)
pulse length results

Conclusions

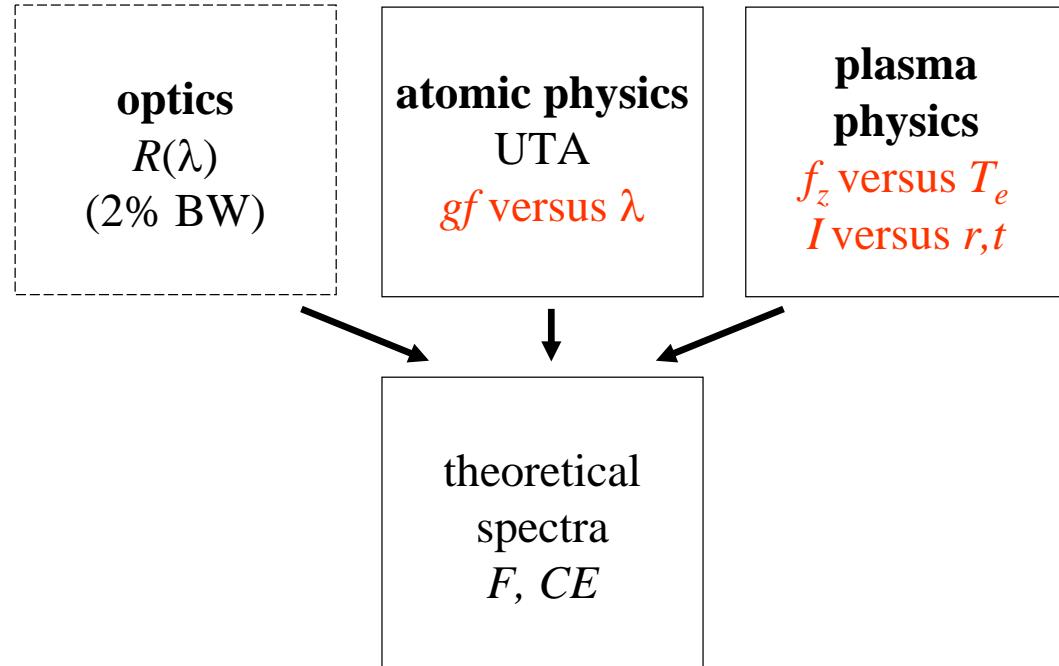
0. Introduction

- quantify 2% in-band emission (13.365-13.635 nm)
- source metric: F (steady-state), CE (time-integrated)

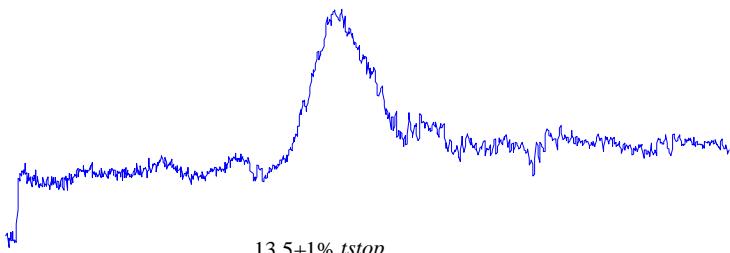
1) “optically thin” – doped tin (5% tin by number)



$$F = \sum_{\lambda} \sum_z R_{\lambda}^n f_z g f_{\lambda,z}$$



2) “optically thick” – pure tin



$$CE = \frac{2\pi \int_{13.5-1\%}^{13.5+1\%} \int_0^{tstop} I_{out}(\lambda, t) dt d\lambda dA}{E_{tot}}$$

experimental comparisons (UCD):
Jenoptik 0.25-m grazing incidence,
1064-nm, 15-ns FWHM Nd:YAG

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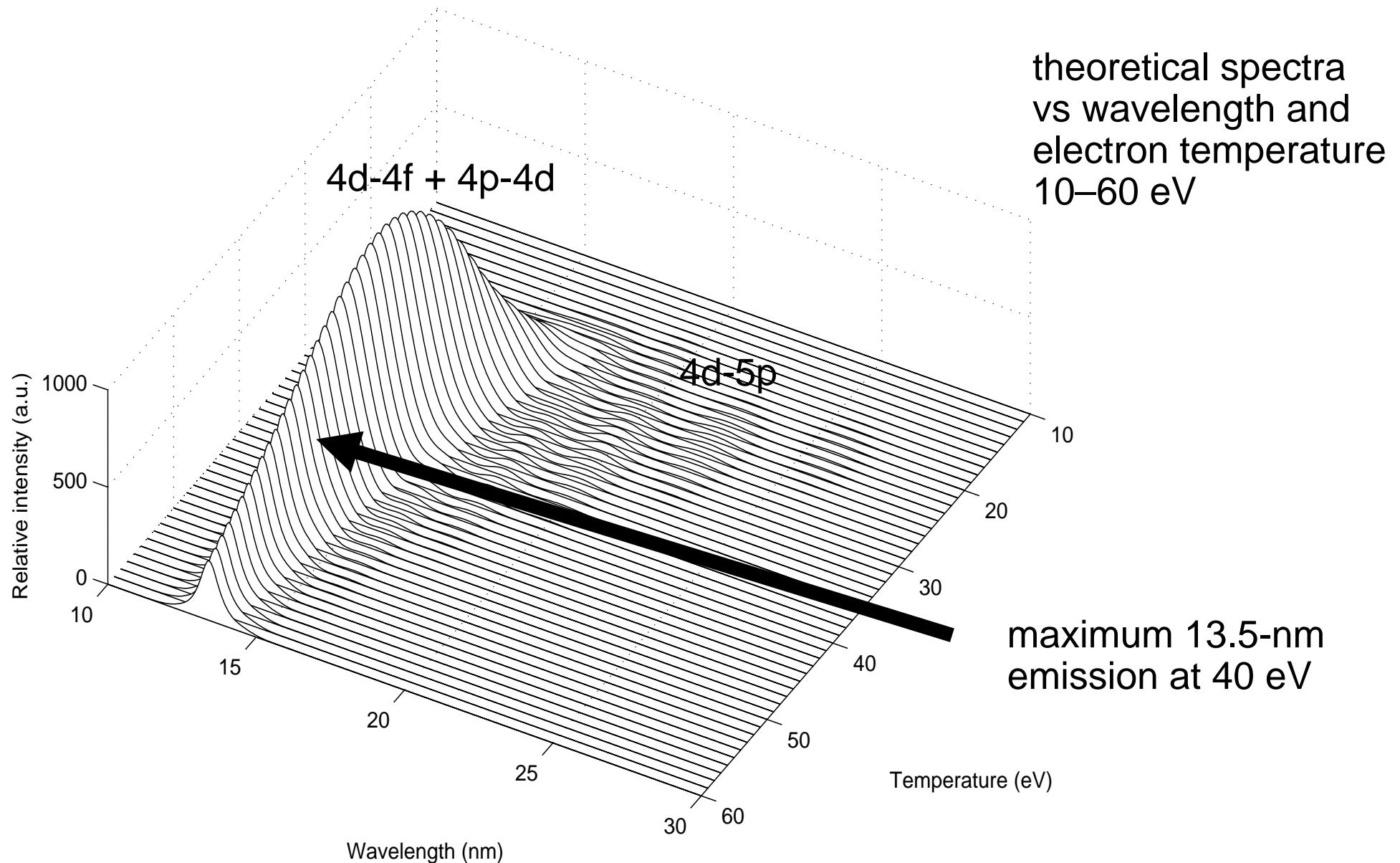
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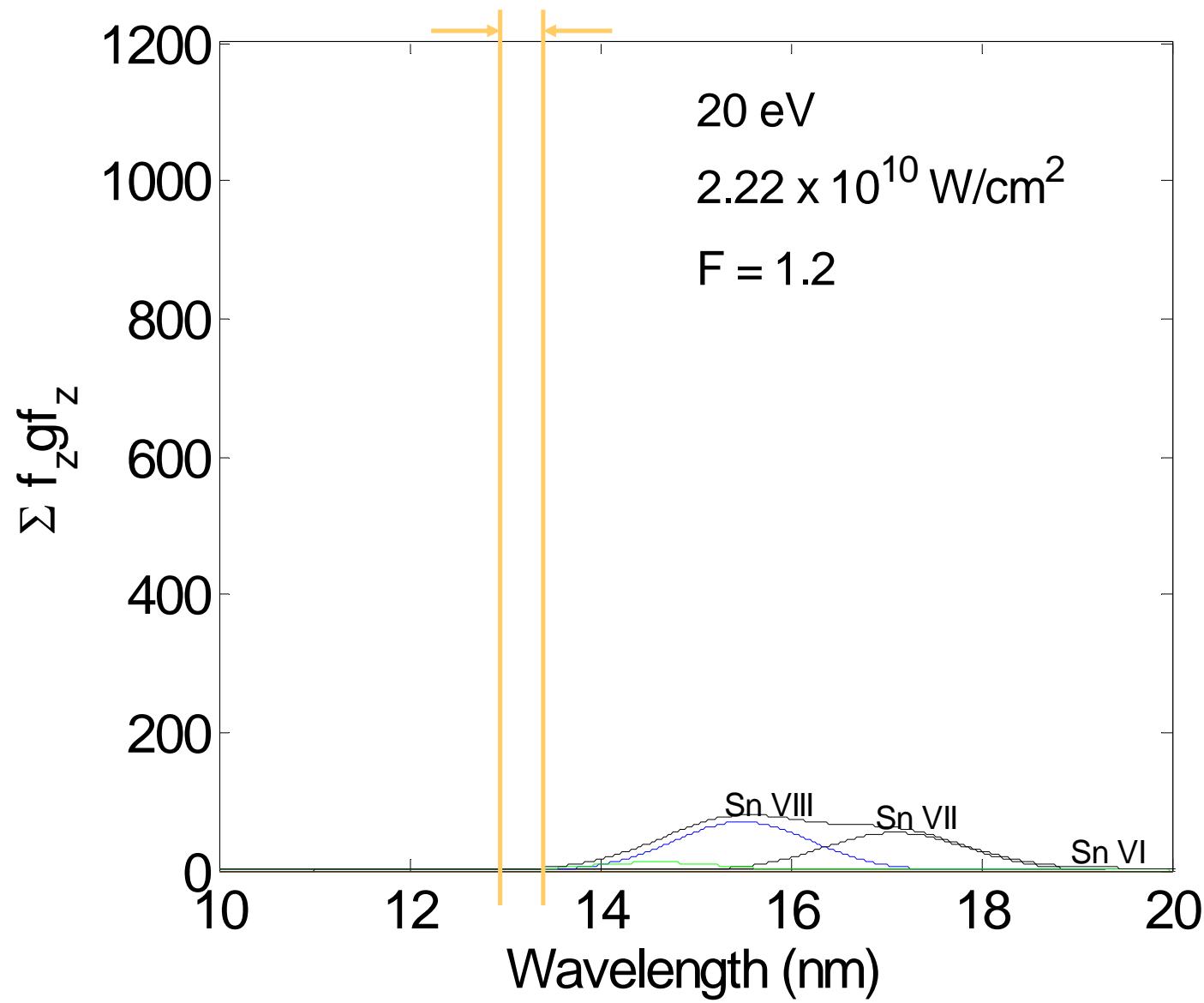
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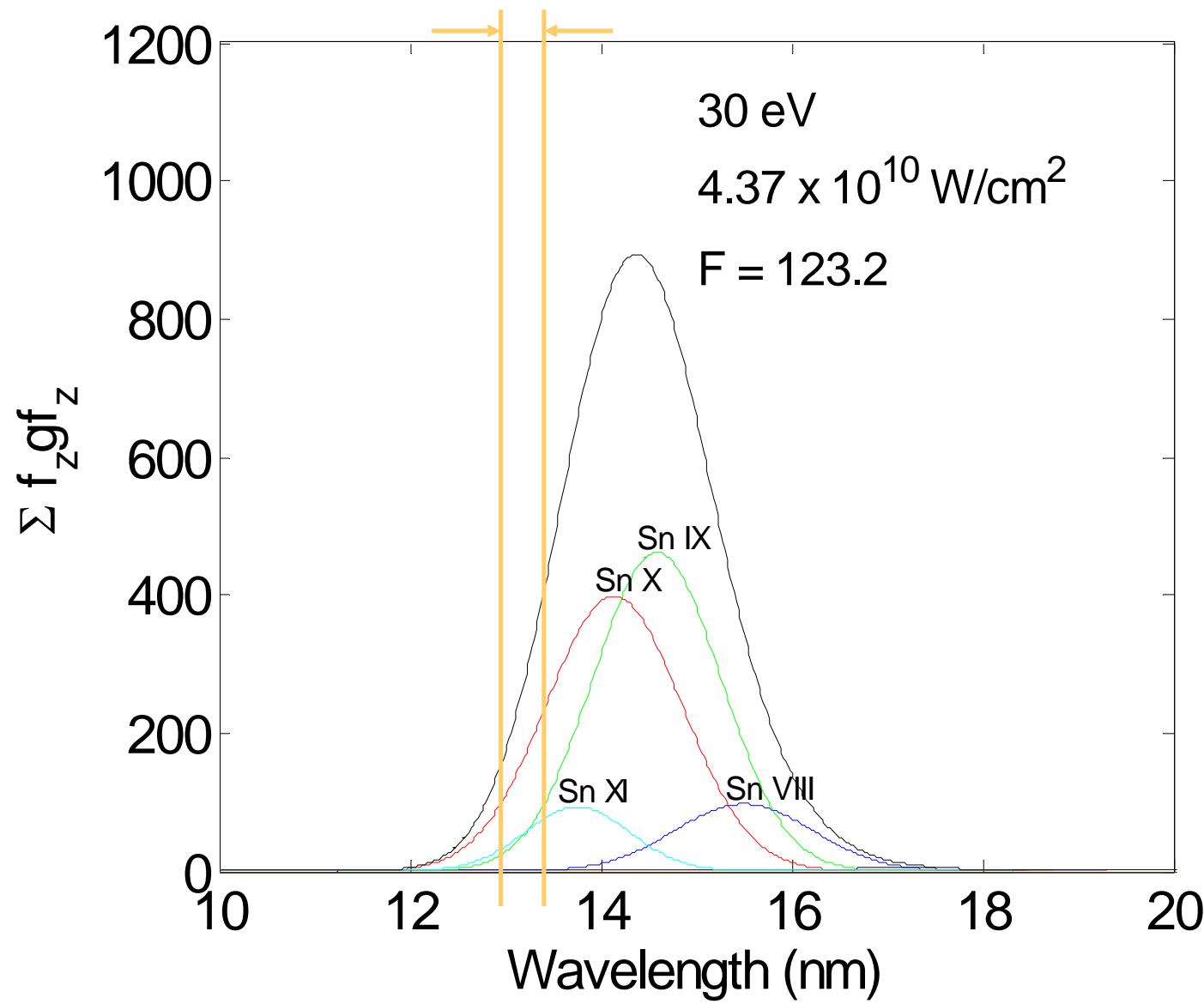
1. ‘Steady-state’ (oscillators weighted by ion distribution) $gf \cdot f_z$



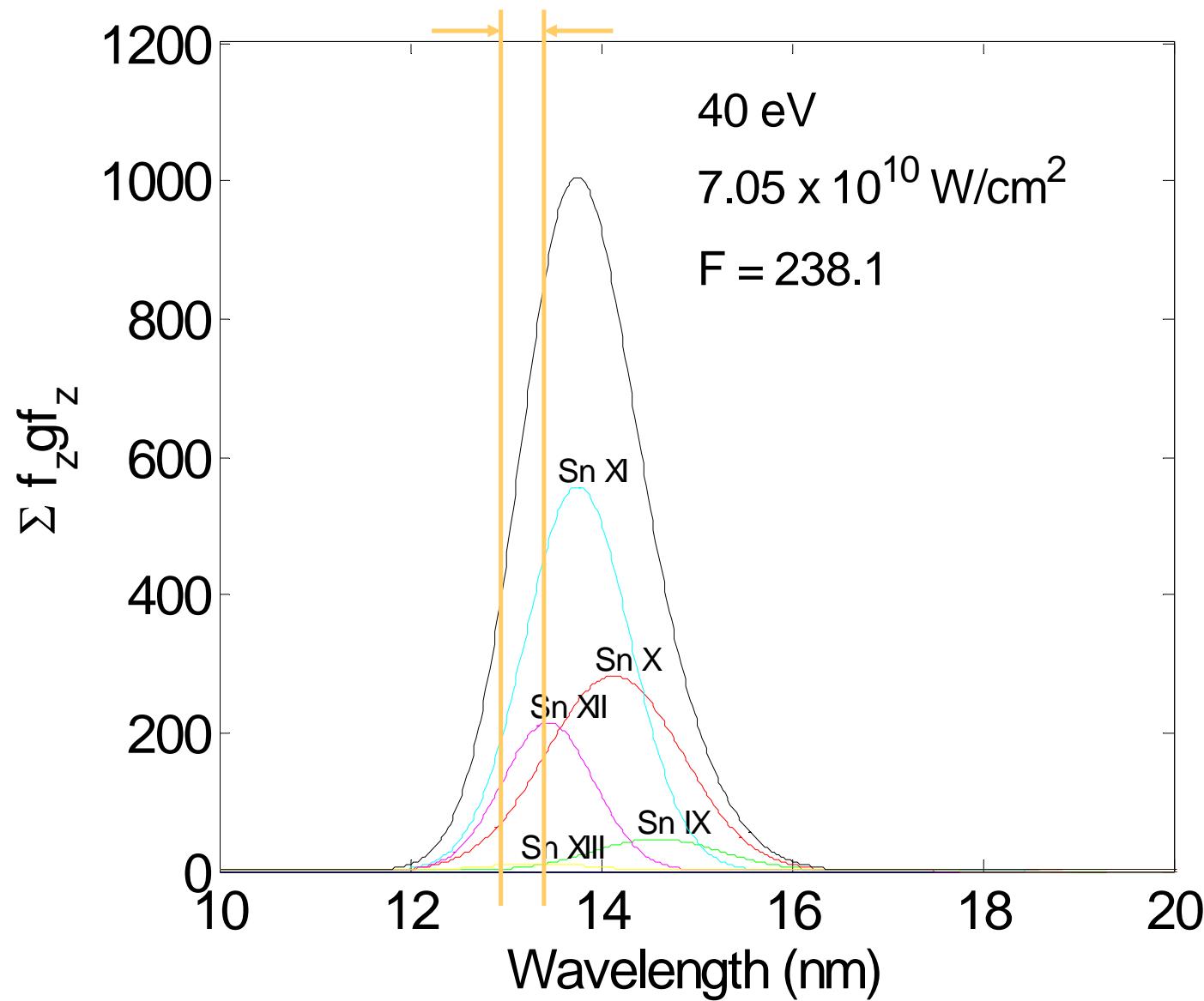
1. Steady-state plasma in-band contribution • 20 eV



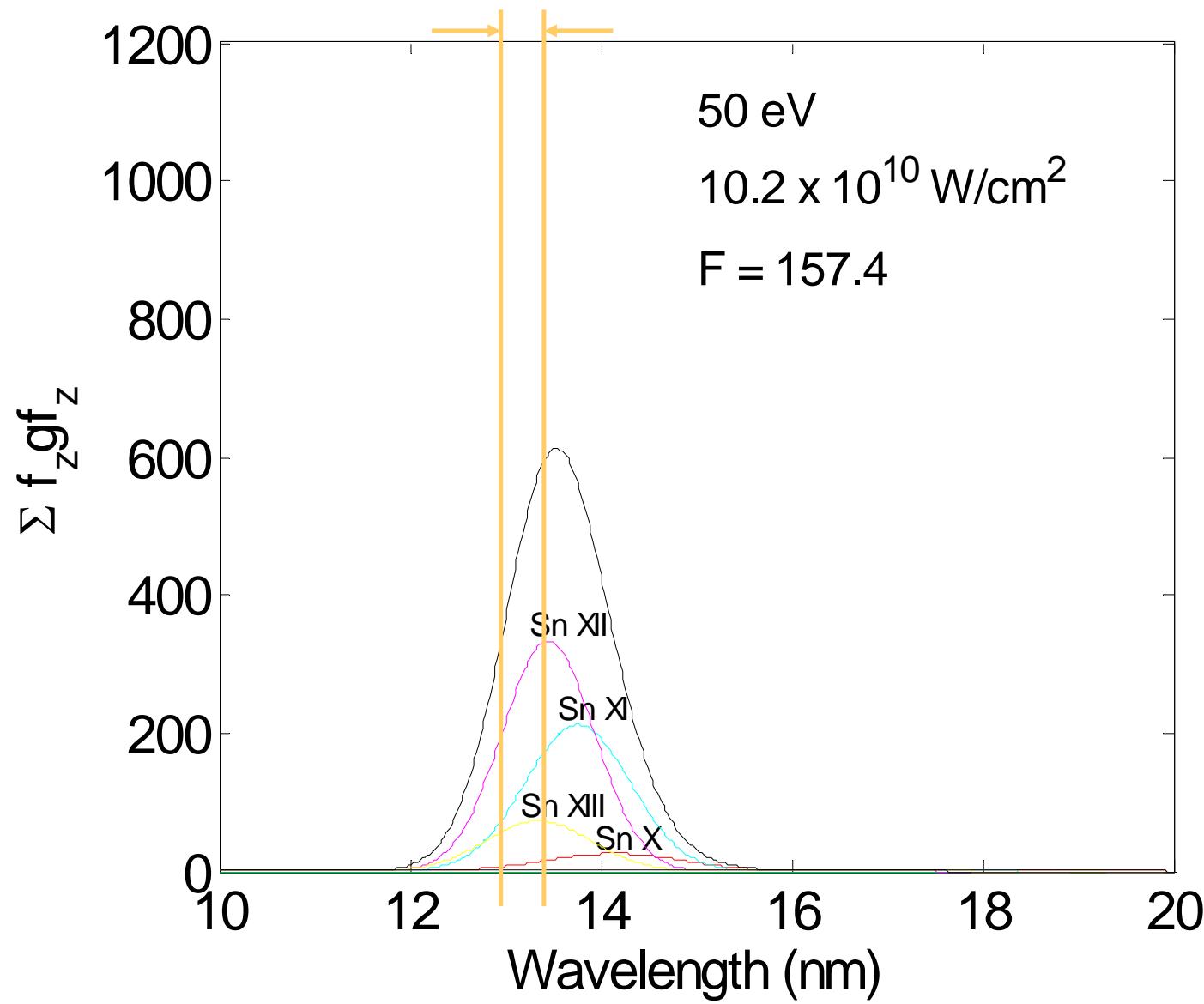
1. Steady-state plasma in-band contribution • 30 eV



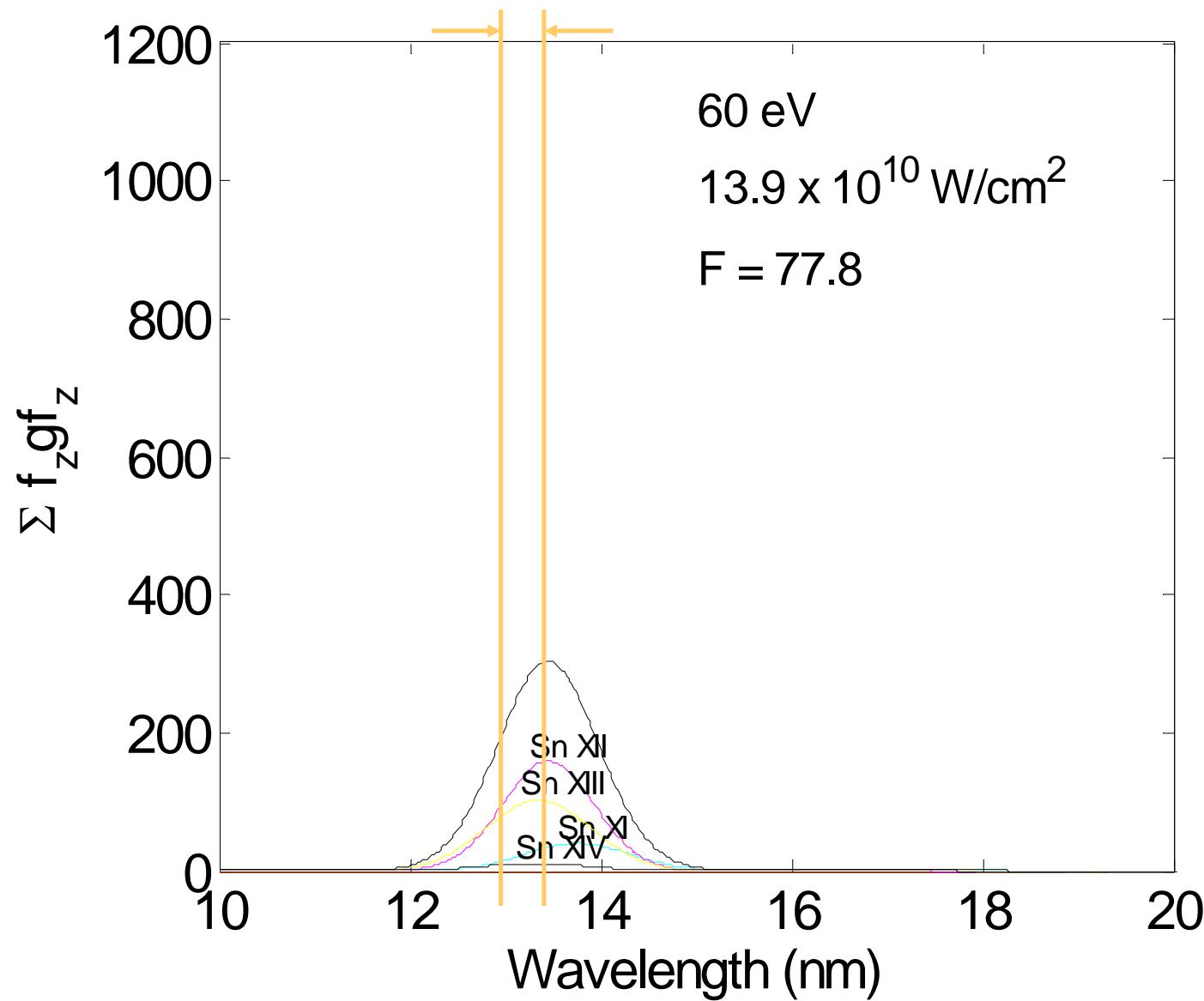
1. Steady-state plasma in-band contribution • 40 eV



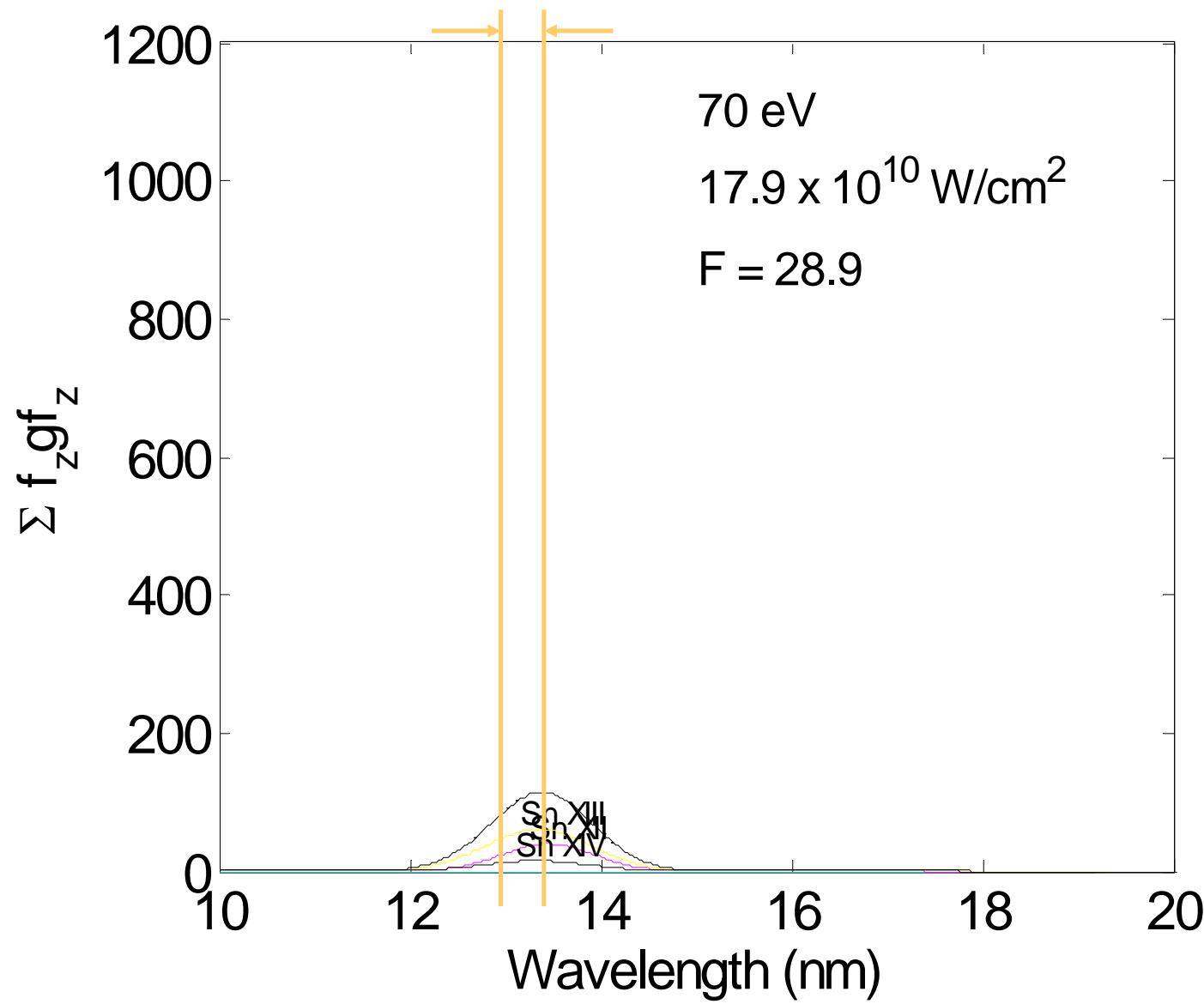
1. Steady-state plasma in-band contribution • 50 eV



1. Steady-state plasma in-band contribution • 60 eV



1. Steady-state plasma in-band contribution • 70 eV

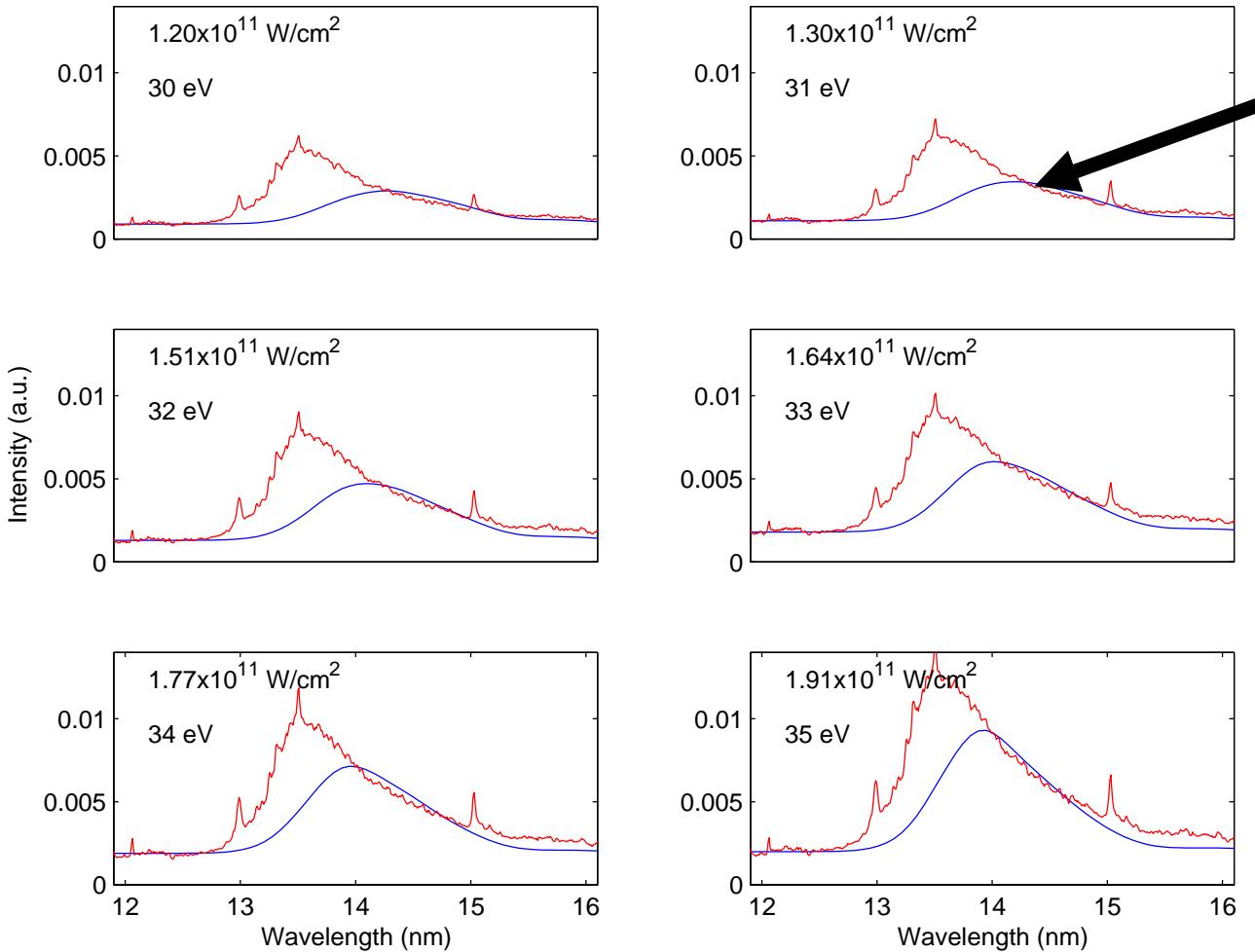


1. Experimental comparison

5% tin by number (optically thin)

Jenoptik 0.25-m grazing incidence flat field spectograph (UCD)

- long wavelength match at lower temperatures
30-35 eV



lower ion stages:
41% Sn IX at 31 eV

J. Appl. Phys. **98** 113301, 2005

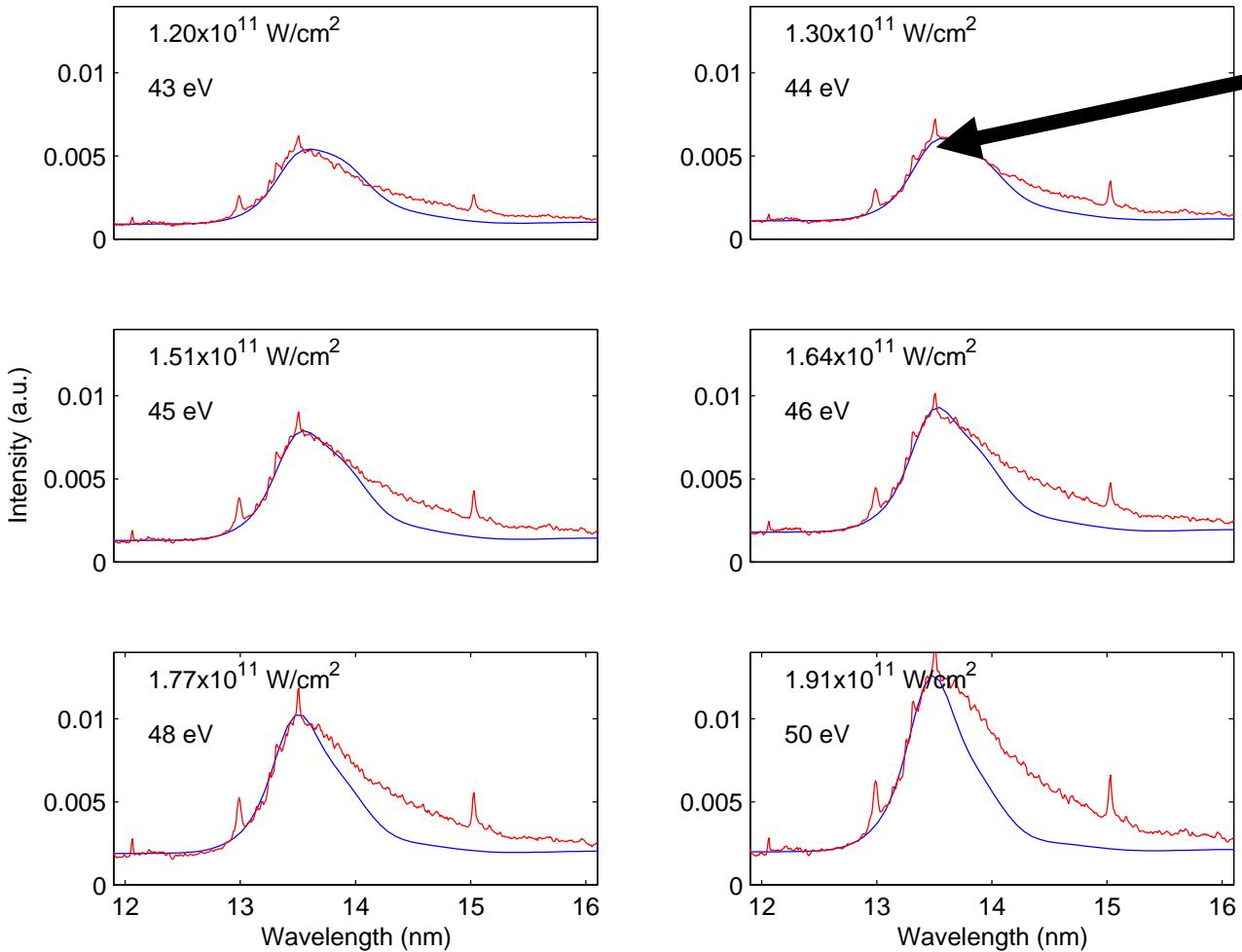
J. Appl. Phys. **99** 093302, 2006

1. Experimental comparison

5% tin by number (optically thin)

Jenoptik 0.25-m grazing incidence flat field spectograph (UCD)

- short wavelength match at higher temperatures
43-50 eV

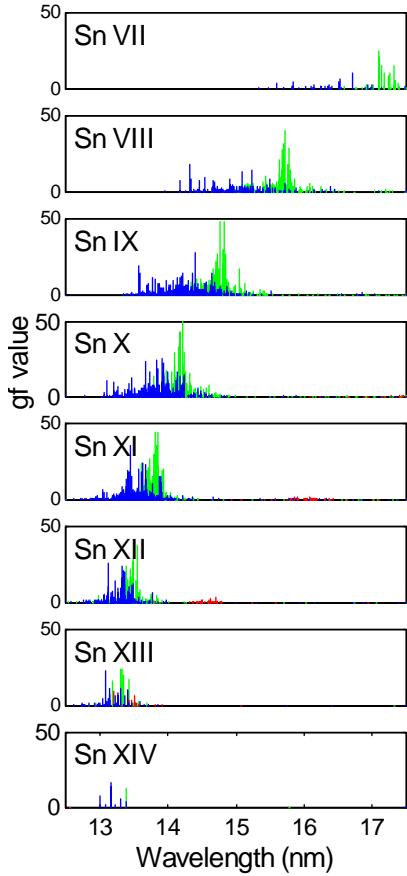


higher ion stages:
43% Sn XII at 44 eV

Note: ground state only
one- T model?

1. Statistical representation '4d' subshell ions: Sn VII–Sn XIV

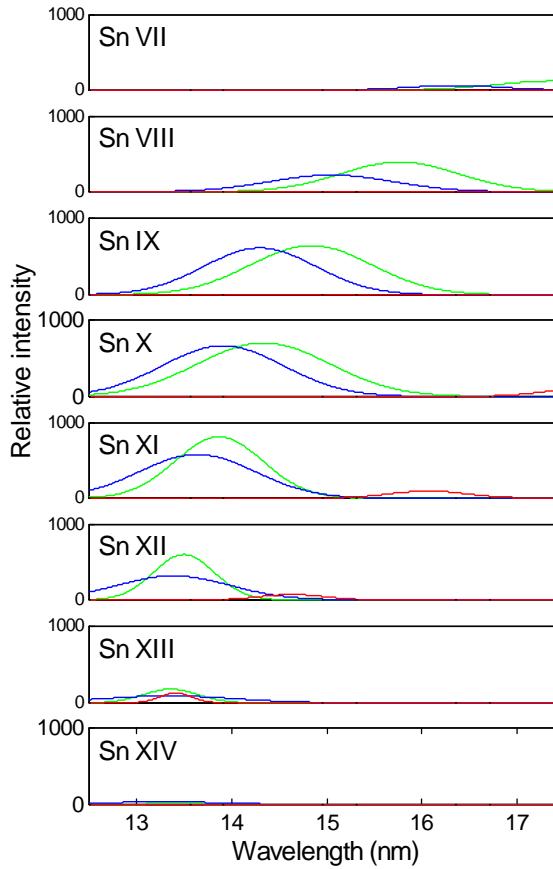
- $\mathbf{H} \psi = E \psi$
- statistical "three-parameter" Gaussian (mean, standard deviation, Σgf) versus λ



$$\mu_n = \frac{\sum_{i=1}^N (\lambda_i)^n gf_i}{\sum_{i=1}^N gf_i}$$



$$f(x) = \frac{1}{\sigma \sqrt{2\pi}} e^{-\frac{(x-\mu_1)^2}{2\sigma^2}}$$



4d → 4f

4p → 4d

4d → 5p

1. Rate equations: ionisation vs recombination

- ion distribution f_z from rate equations
(from atomic processes)
 - collisional ionisation, $S(z-1)$
 - radiative recombination, $\alpha_r(z)$
 - three-body recombination, $\alpha_{3b}(z)$
 - semi-empirical
- a balance of ionisation and recombination

$$f_z = \frac{S}{\alpha_r + n_e \times \alpha_{3b}}$$

ionisation

 S

recombination

$$S = 9 \times 10^{-6} \xi_z (T_e / \chi_z)^{1/2} e^{(-\chi_z / T_e)} \chi_z^{3/2} (4.88 + T_e / \chi_z)$$

$$\alpha_r = 5.2 \times 10^{-14} (\chi_z / T_e)^{1/2} Z [0.429 + .5 \log(\chi_z / T_e) + 0.469 (T_e / \chi_z)^{1/2}]$$

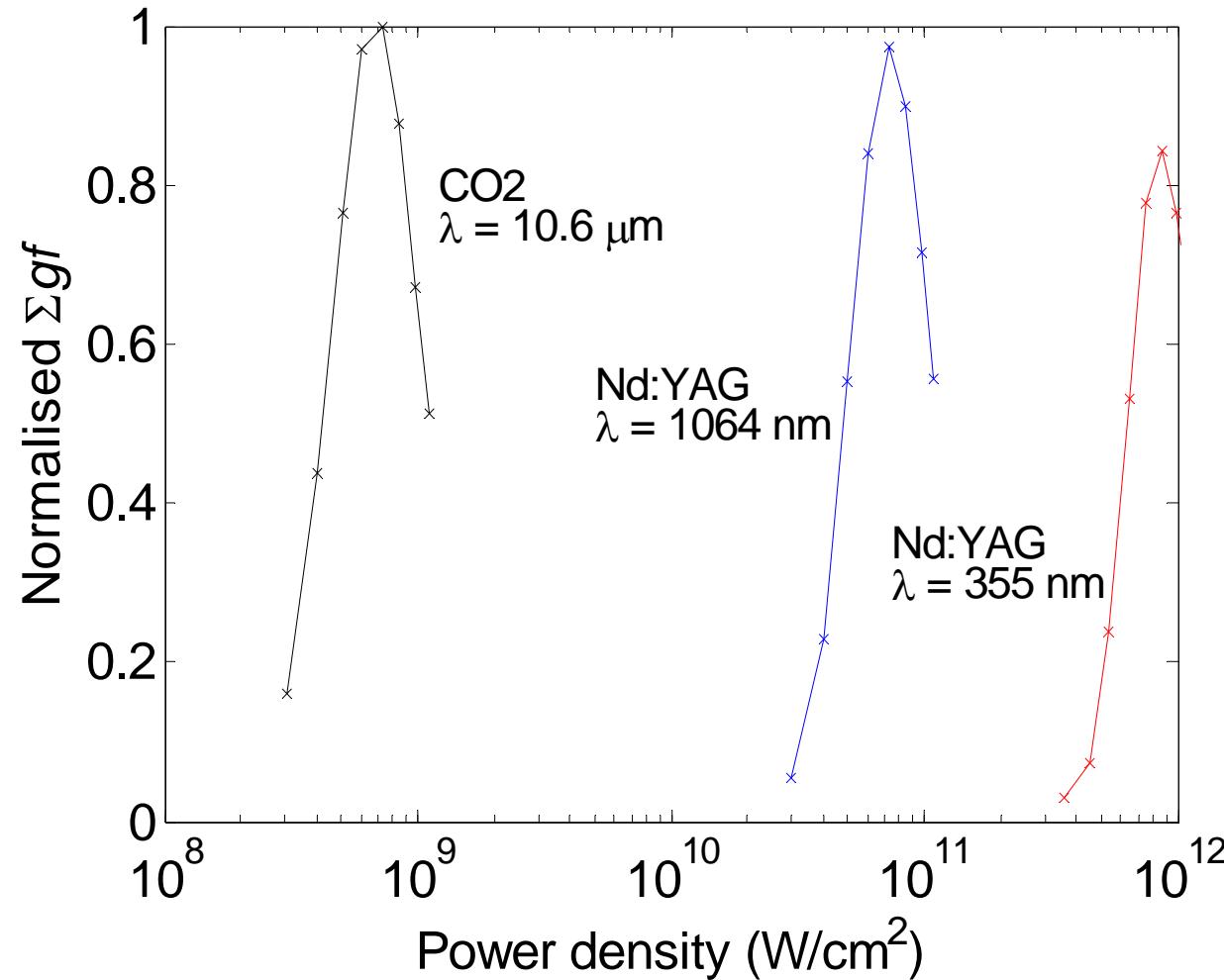
$$\alpha_{3b} = 2.97 \times 10^{-27} \xi_z T_e \chi_z^2 (4.88 + T_e / \chi_z)$$

S = collisional ionisation,
 α_r = radiative recombination,
 D_j = dielectronic recombination,
 α_{3b} = three-body recombination,
 T_e = electron temperature,
 n_e = electron density,
Z = atomic number,
 χ_z = ionisation potential,
 ξ_z = number of open shell electrons,
z = ion stage.

1. Nd:YAG and CO₂

Optimising 13.5-nm emission as a function of laser wavelength

- $n_e(\text{CO}_2) = 1/100 n_e(\text{YAG})$
- lower power density to achieve 40 eV ($\lambda^2\phi = \text{constant}$)



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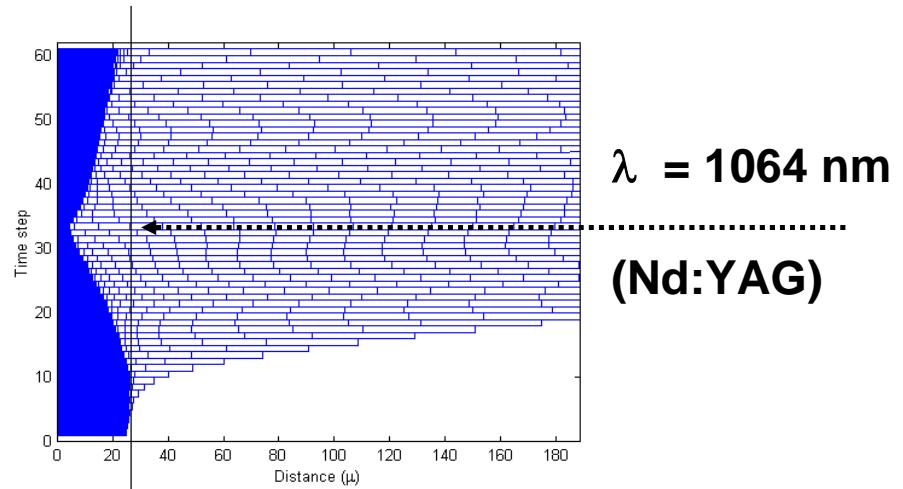
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2. What is the emission as a function of time/space?

time-dependent plasma
 $dn/dt \neq 0$



time-dependent: sharp temperature gradients

$t(\text{atomic processes}) \sim t(\text{laser})$

n_e and $T_e = f(r, t)$

c.f. steady-state: one temperature
 $t(\text{atomic processes}) < t(\text{laser})$

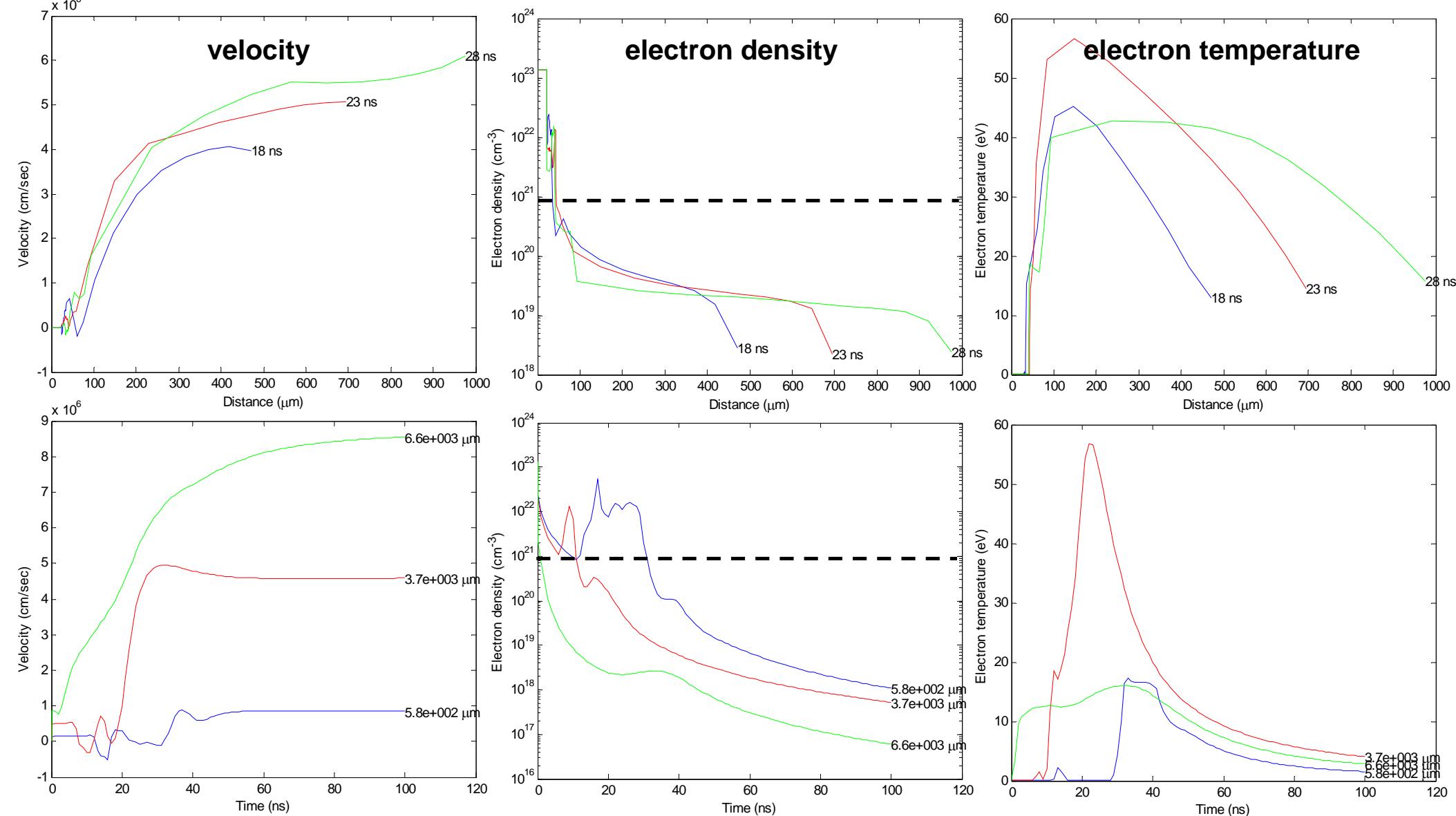
$n_e = n_{ec} \approx 10^{21} \lambda^{-2}$ and $T_e \propto Z^{1/5} (\lambda^2 \phi)^{3/5}$

- FORTRAN laser-plasma interaction code Medusa
- one-dimensional **Lagrangian** difference mesh
- hydrodynamics variables: **density** (ρ), **velocity** (v), **ion temperature** (T_i), **electron temperature** (T_e) as functions of **space** (r) and **time** (t) from Navier-Stokes equations
- **average atom model** (to simplify rate equations especially for high Z). $P_n = n$ shell occupation number.
- excitation/de-excitation included (I -degenerate level populations calculated from average atom model)

2. Sn hydrodynamics: vs (r,t)

versus cell at: 18, 23, 28 ns (POP: 23 ns) 15-ns FWHM

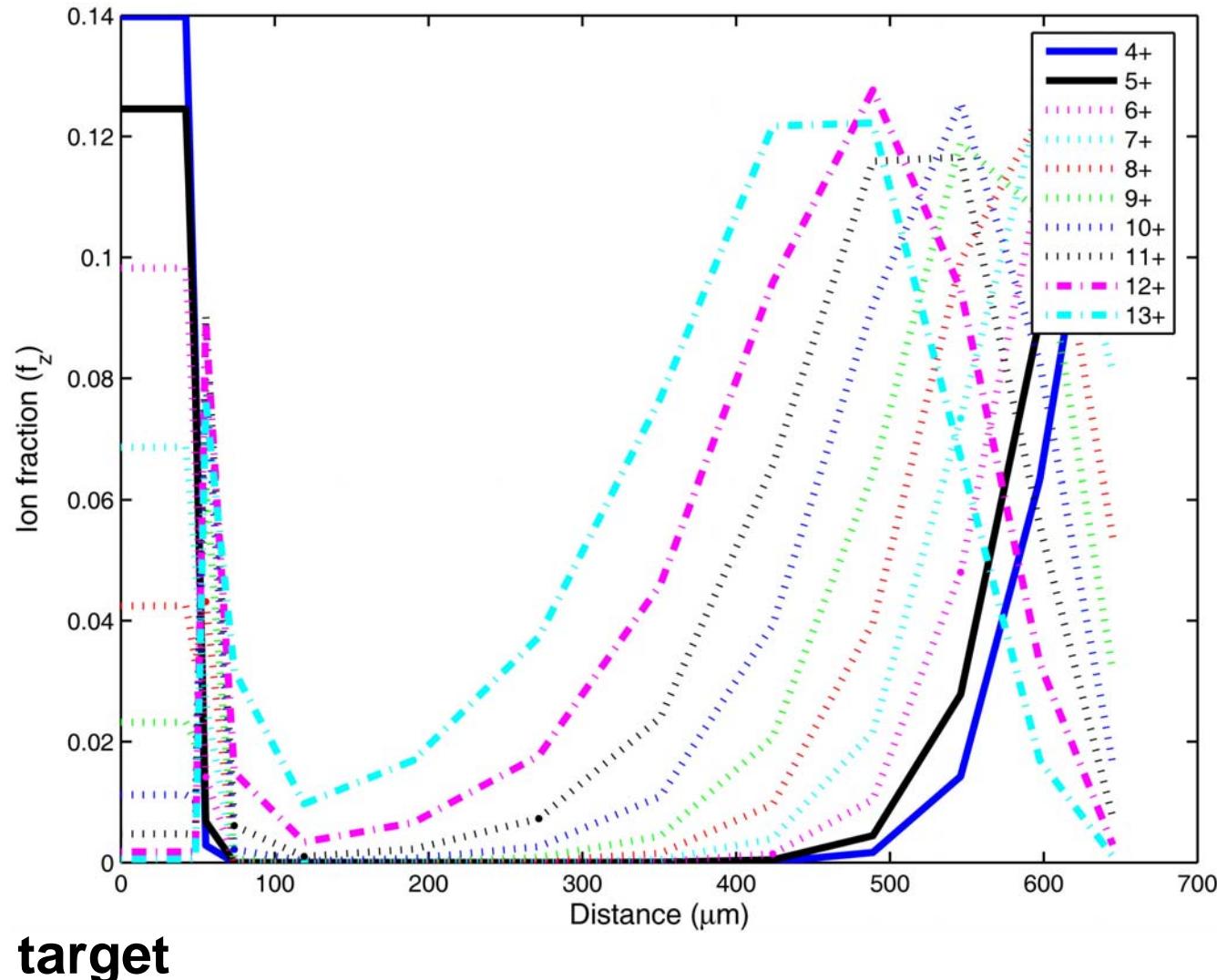
laser: $P_{\text{equiv}} = 1.3 \times 10^{11} \text{ W/cm}^2$
 target: solid tin, 50 μm -diameter wire
 simulation: 400 cells, 100 ns, cylindrical geometry



versus time at: 400, 392, 384

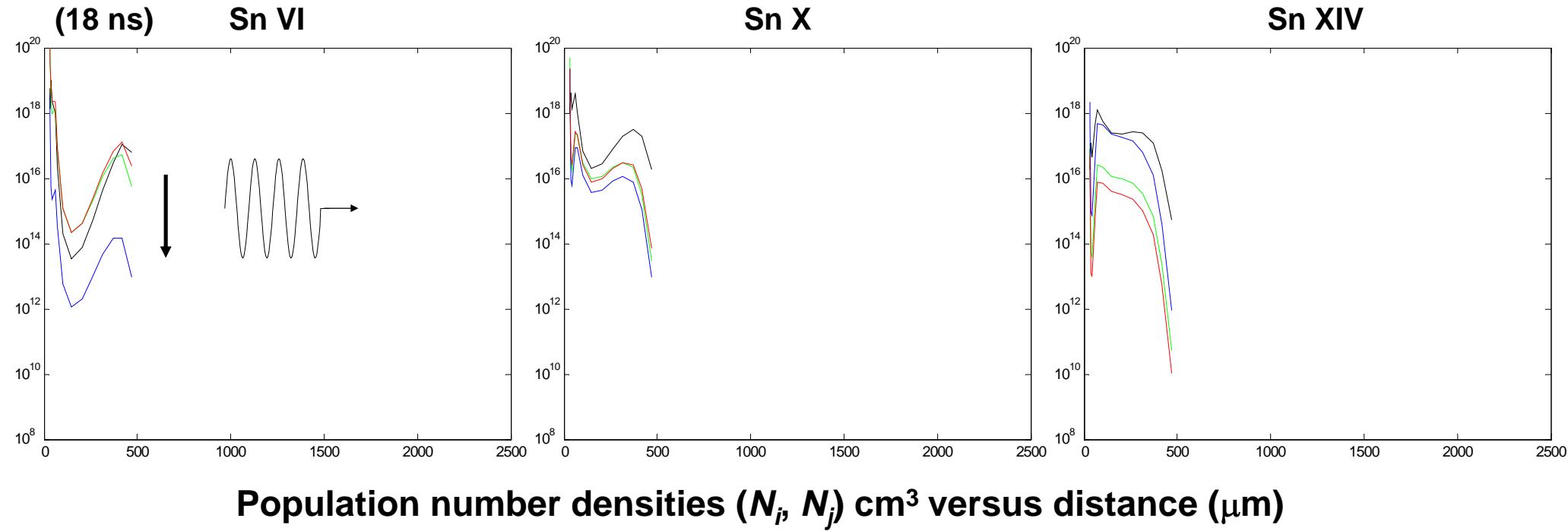
2. Ion distribution = $f(r,t)$ 1-D x time Peak of the pulse

- Sn^{4+} -- Sn^{13+} ion fraction and ion density (cm^{-3})



2. Time-dependent plasma: Sn level populations (for $\Delta n = 0$)

- Sn V – Sn XIV UTA (/ degeneracy removed with energy functional)

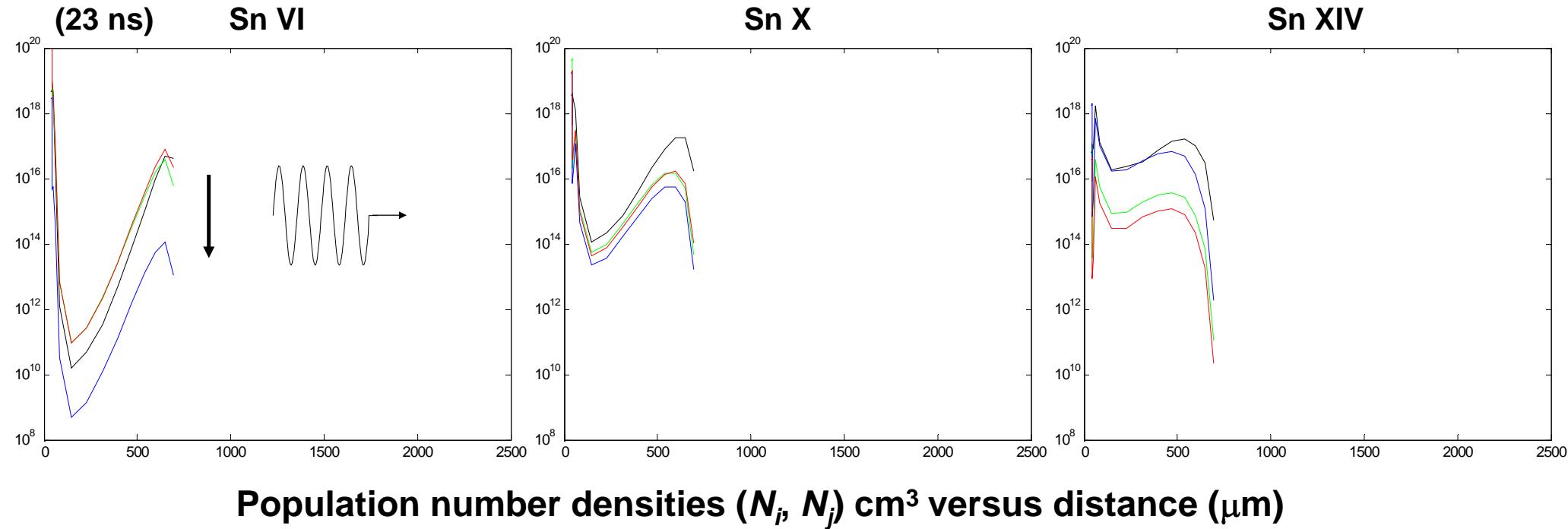


E_0 : black: $4p^64d^N$
green: $4p^64d^{N-1}4f^1$
blue: $4p^54d^{N+1}$
red: $4p^64d^{N-1}5p^1$

J. Appl. Phys. **101** 043301, 2007

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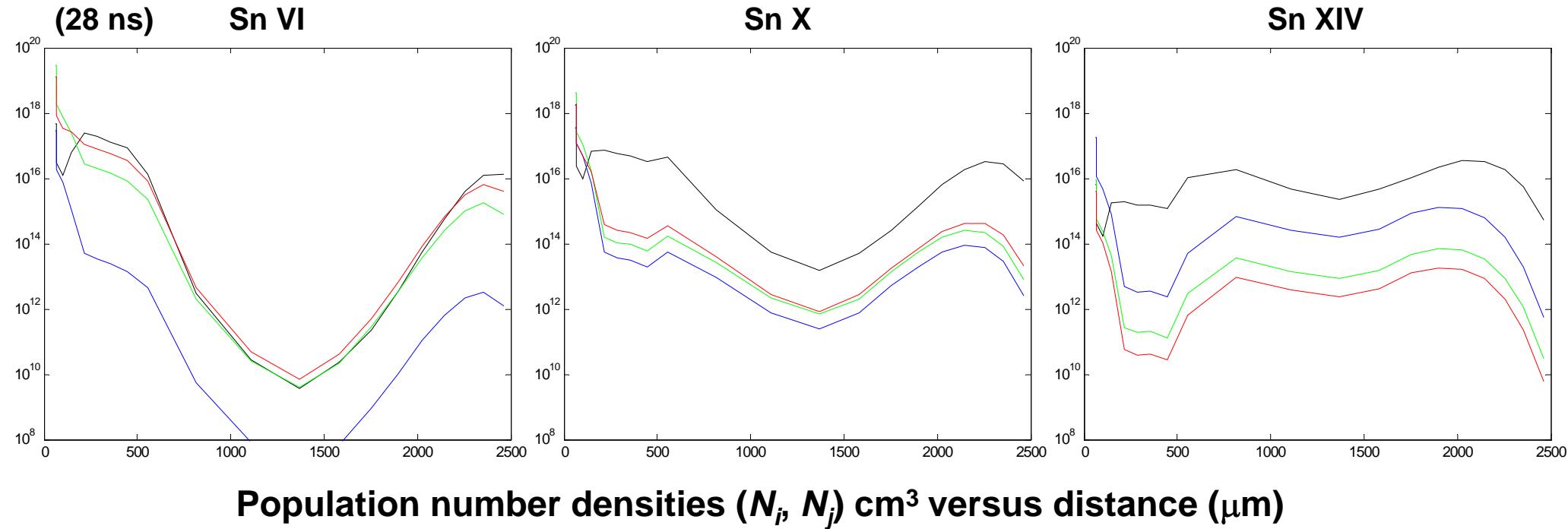


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2. All together: sum up intensity with absorption

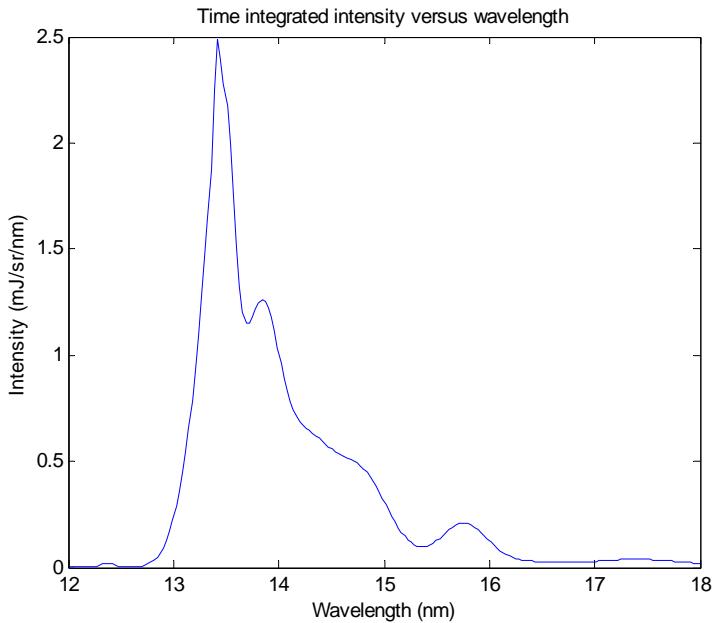
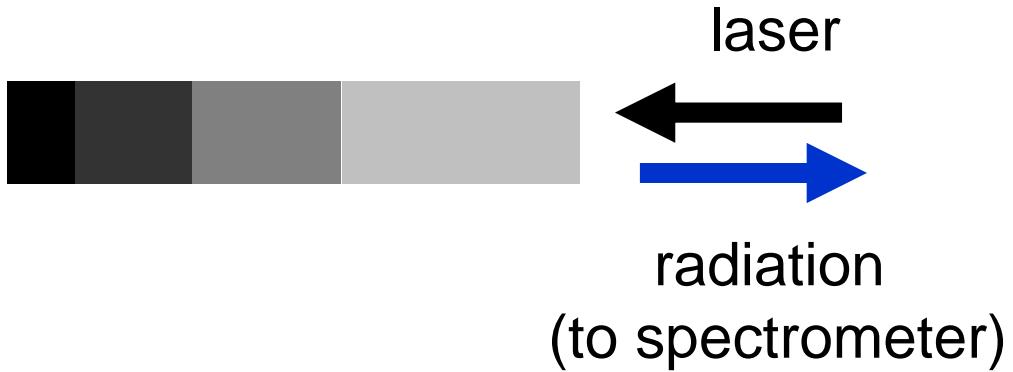
- $I_m = f(r, t) \rightarrow I = f(\lambda, t)$
- absorption strength, f_{ij} (Cowan, stat UTA) $f_{ij} = \frac{gf}{D_i}$
- level populations, N_{ij} (Medusa, E. F.) $N_{ij} = W_a \frac{\rho}{Am_p}$
- degeneracy, g or D
- cell width, Δz (Medusa)
- absorption cross section, χ

$$I_m^{obs}(\nu) = I_m(\nu) \exp\left(\sum -\chi_n(\nu) \Delta z_n\right)$$

$$\chi_m^l(\nu) = \frac{\pi e^2}{4\pi \epsilon_0 m_e c} f_{ij} N_i \left[1 - \frac{N_j g_i}{N_i g_j} \right] \phi_m^l(\nu)$$

$$I_m(\nu) = \sum_l S_m^l(\nu) [1 - \exp(\chi_m^l(\nu) \Delta z_m)]$$

$$S_m^l(\nu) = \frac{2h\nu^3}{c^2} \frac{1}{\left[\frac{N_i g_j}{N_j g_i} - 1 \right]}$$



The emission as a function of time, I_{out} , is calculated by summing up attenuated cell emission over all m cells

n/l level populations as a function of r,t

for $4p^6 4d^N$, $4p^6 4d^{N-1} 4f^1$, $4p^5 4d^N$, and $4p^6 4d^{N-1} 5p^1$ configurations

The rest is atomic accounting (a.k.a. computational physics)

10 ions

3 transitions

6 shells (SHM)

400 x 400 cells (summed attenuated absorption)

or 200 x 401 = 80200 cells

100 timesteps

200 frequency points in 2% bandwidth

$\sim 3 \times 10^{11}$ (300 billion) calculations/spectrum

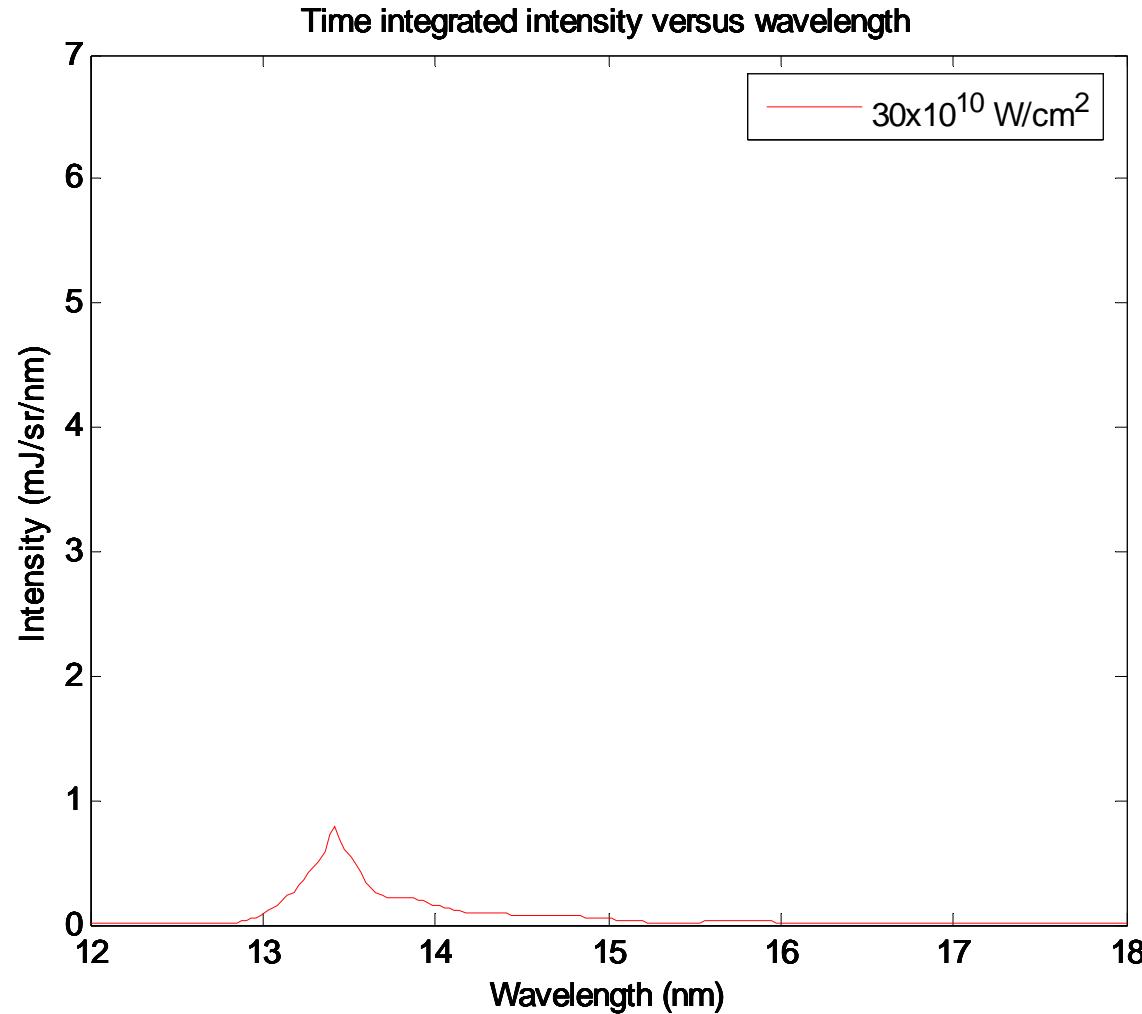
But the good news: **only 3 UTA transitions instead of 100,000 in the line-by-line analysis**

So 300 billion calculations/spectrum instead of 10 thousand trillion!

2. Sn time-integrated UTA spectra

- power density from 0.5 to $3.0 \times 10^{11} \text{ W/cm}^2$

laser: 1.064 μm , 15.0 ns (FWHM), Gaussian
target: solid tin, 90 μm -diameter wire
simulation: 400 cells, 100 ns, cylindrical geometry

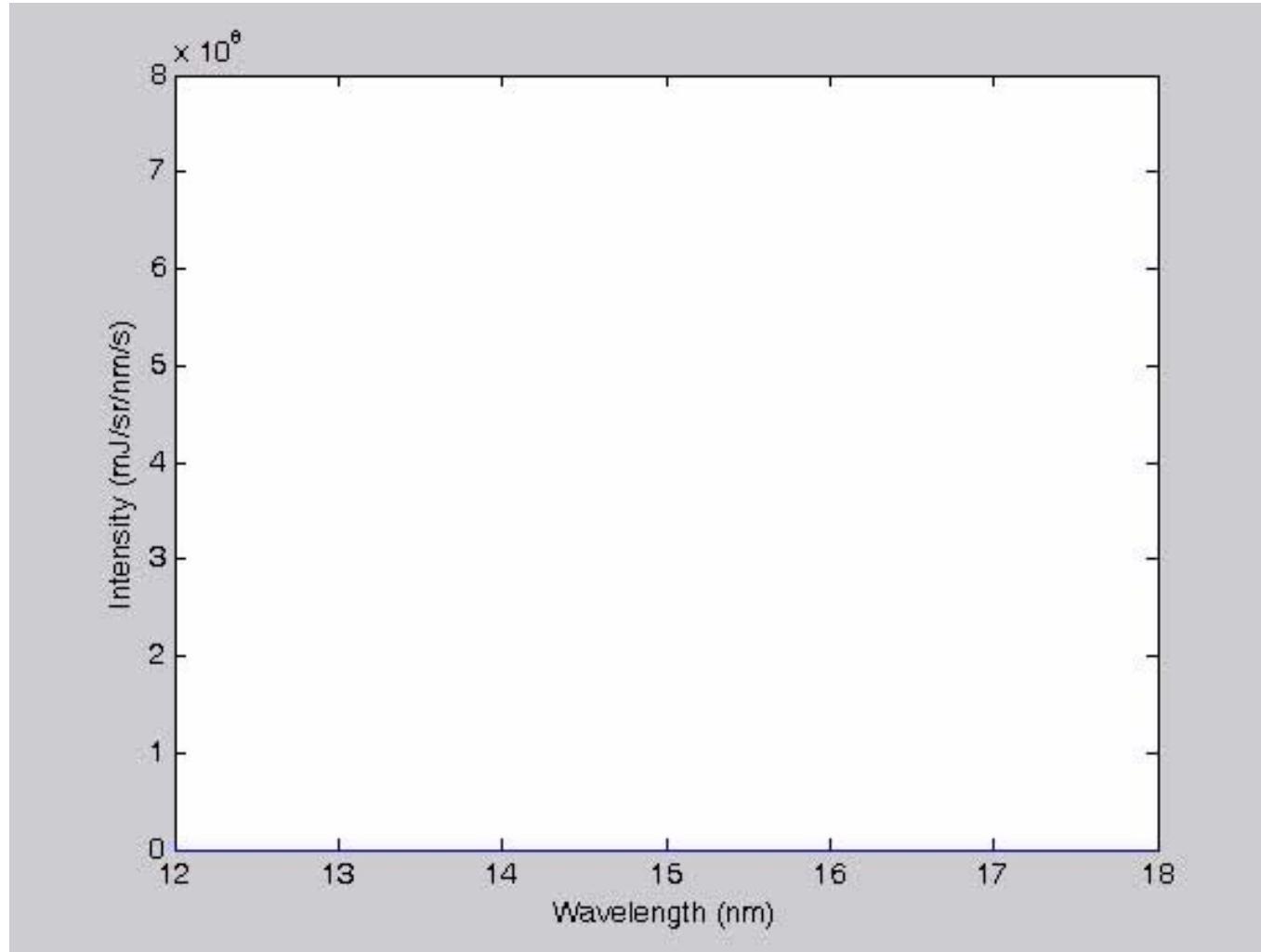


Calculated spectra: intensity versus wavelength

2. Sn UTA spectra versus time

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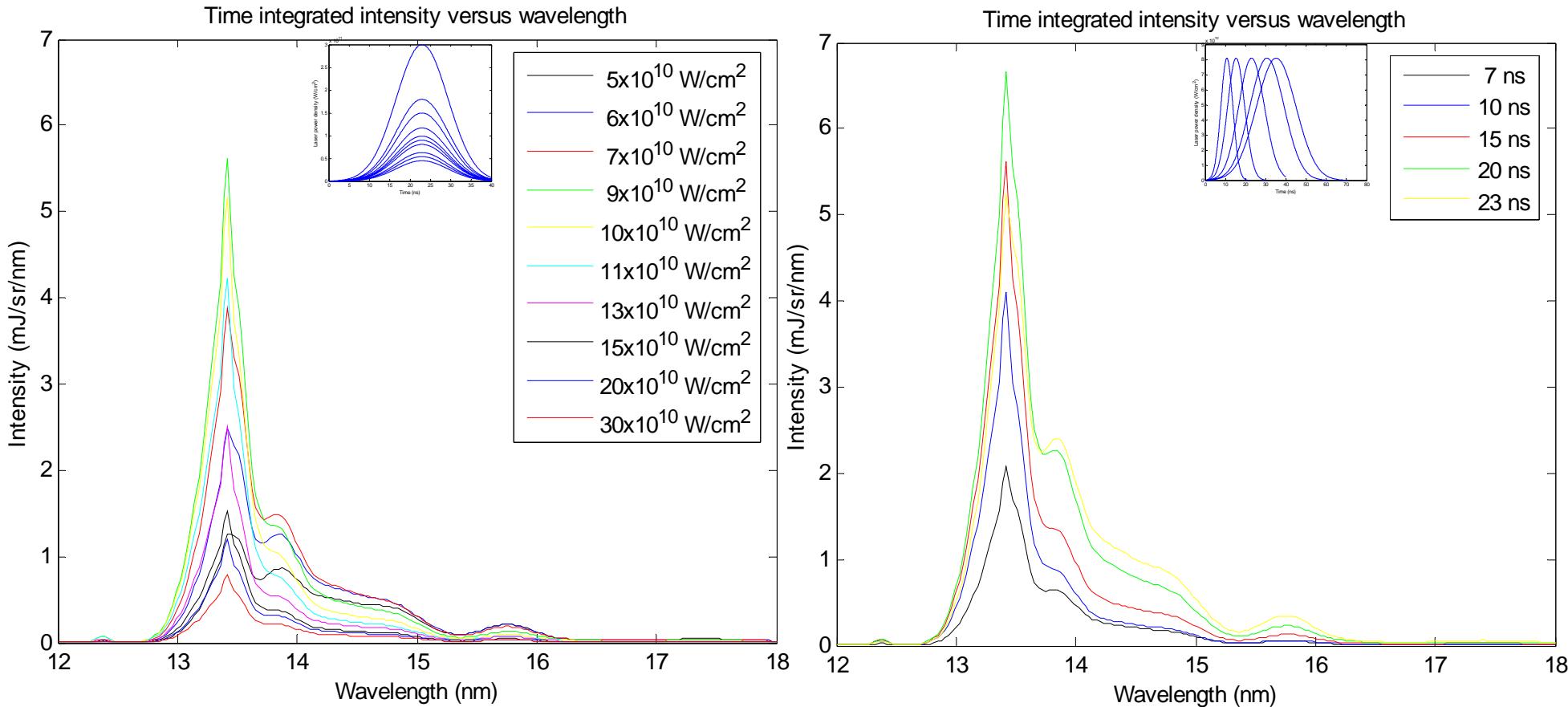
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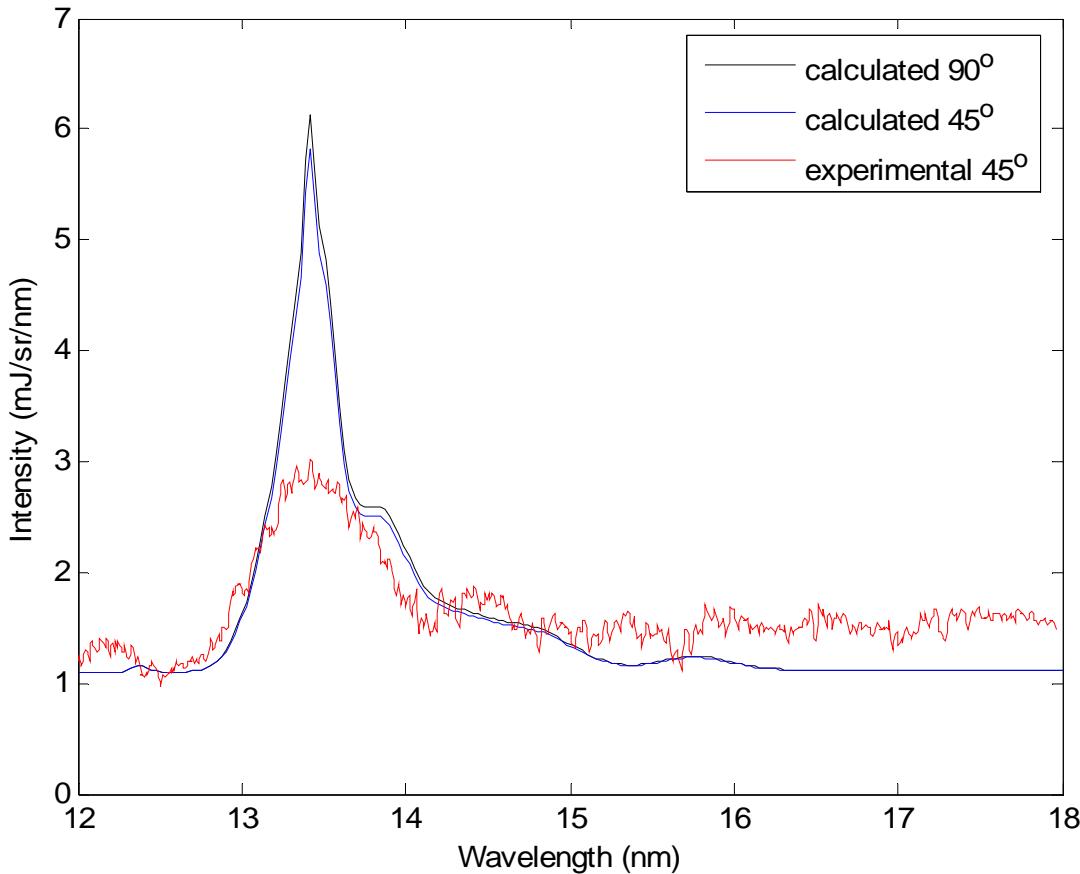
- power density from 0.5 to $3.0 \times 10^{11} \text{ W/cm}^2$: **max CE at $0.8 \times 10^{11} \text{ W/cm}^2$**
- pulse width from 7 to 23 ns: **max CE for 10-ns FWHM (4% > than at 15-ns)**



Calculated spectra: intensity versus wavelength

2. Experimental comparison

- Calculated: $0.8 \times 10^{11} \text{ W/cm}^2$
- Experimental: $0.9 \times 10^{11} \text{ W/cm}^2$



Experimental and calculated spectra

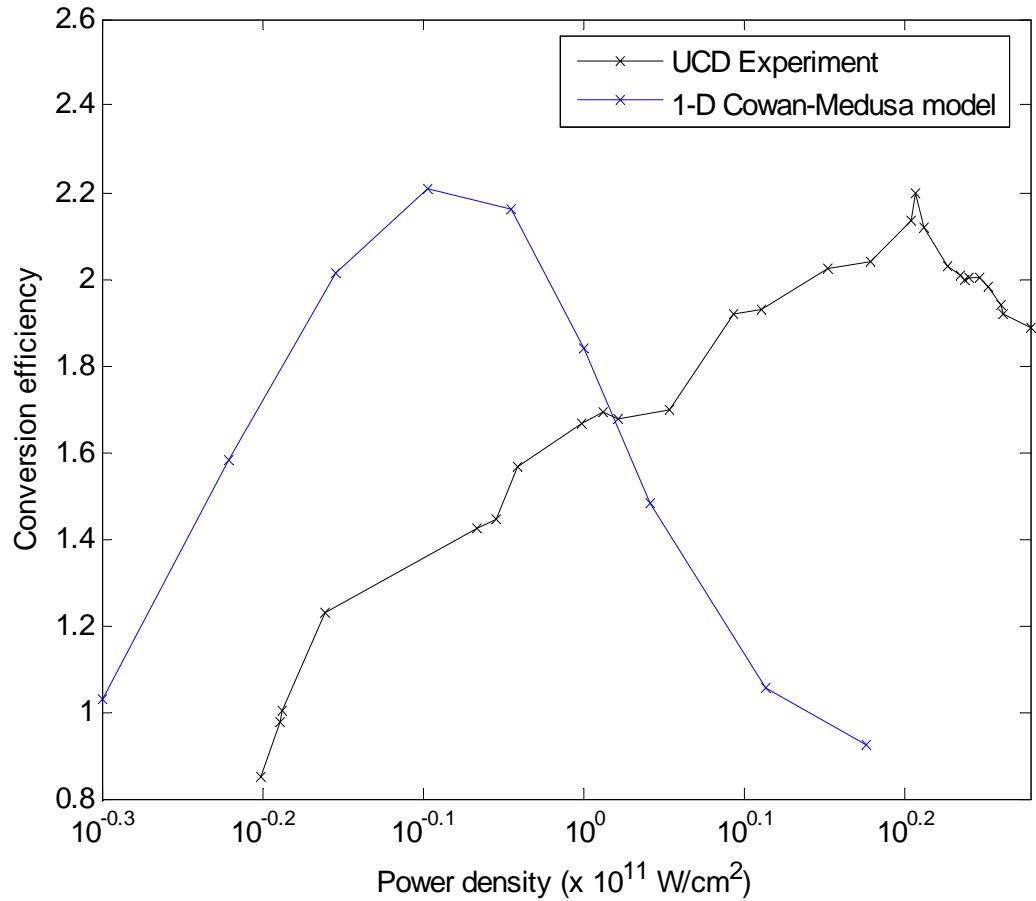
1-D plasma

- ✓ multi electron temperature
- ✓ multi electron density
- ✓ energy level populations

J. Appl. Phys. **99** 093302, 2006

2. Conversion efficiency

- power density survey: max CE at $0.8 \times 10^{11} \text{ W/cm}^2$
- Lateral expansion not included!



$$CE = \frac{2\pi \int_{13.5-1\%}^{13.5+1\%} \int_0^{t_{stop}} I_{out}(\lambda, t) dt d\lambda dA}{E_{tot}}$$

Conversion efficiency (%) versus power density

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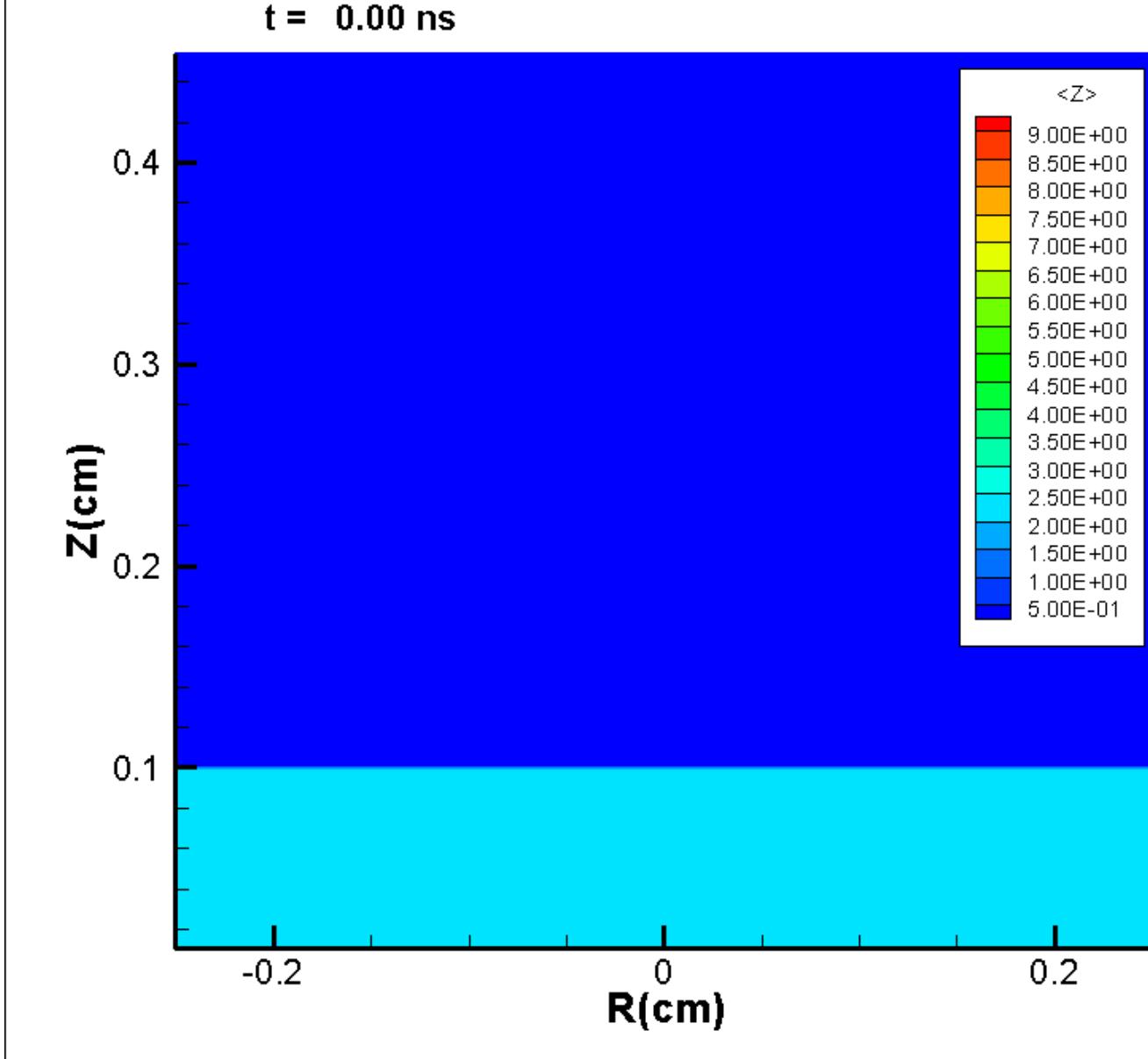
Conclusions

3. 2D plasma

Z* EPPRA

Implicit, E-L MHD,
average atom code,
solves ionisation
kinetics self-
consistently with
radiation transfer

82 x 69 grid, YAG

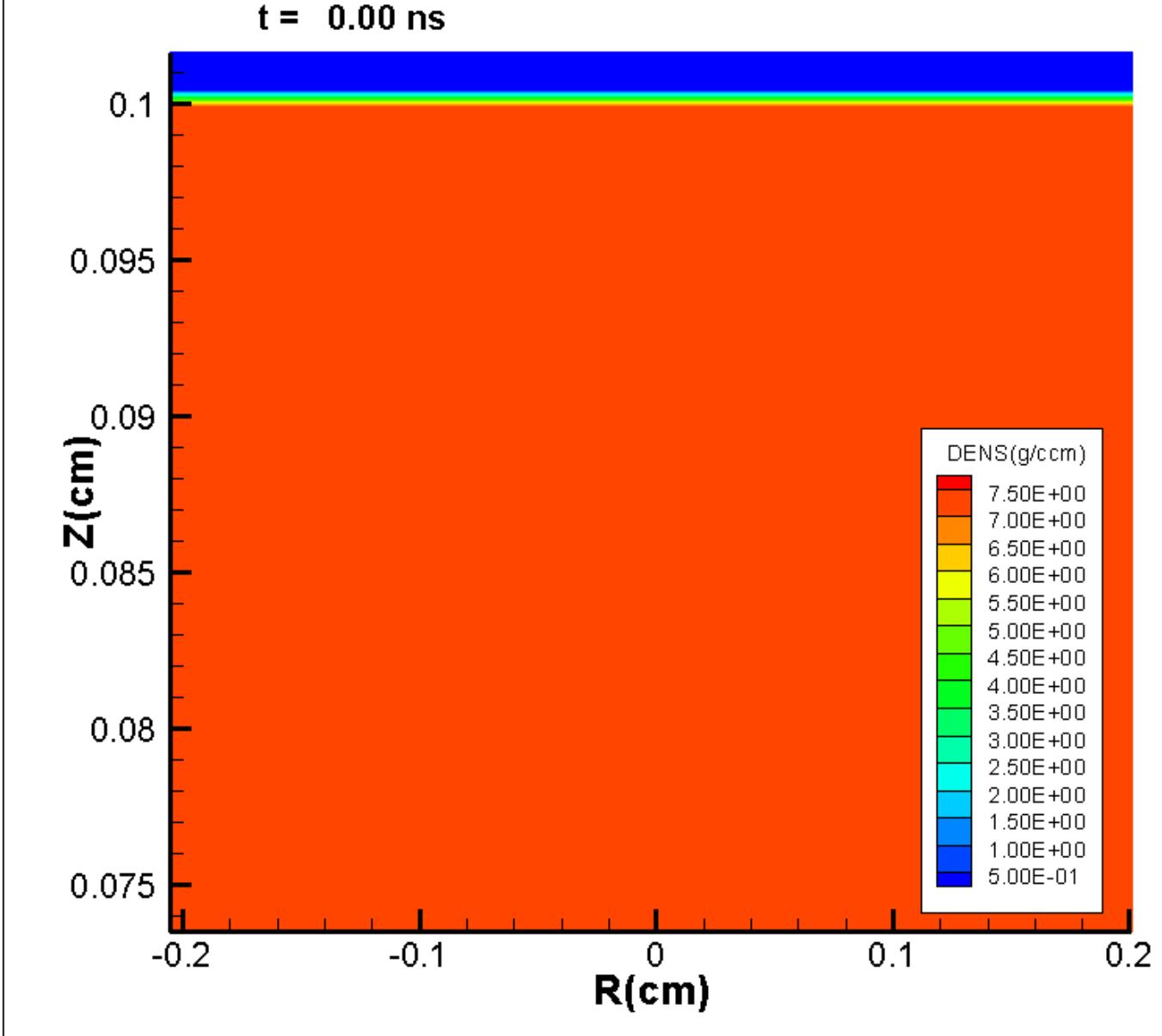


3. 2D plasma

Z* EPPRA

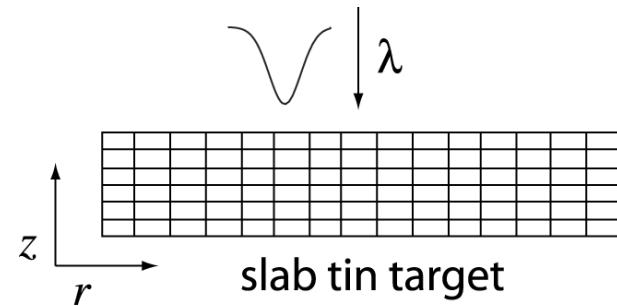
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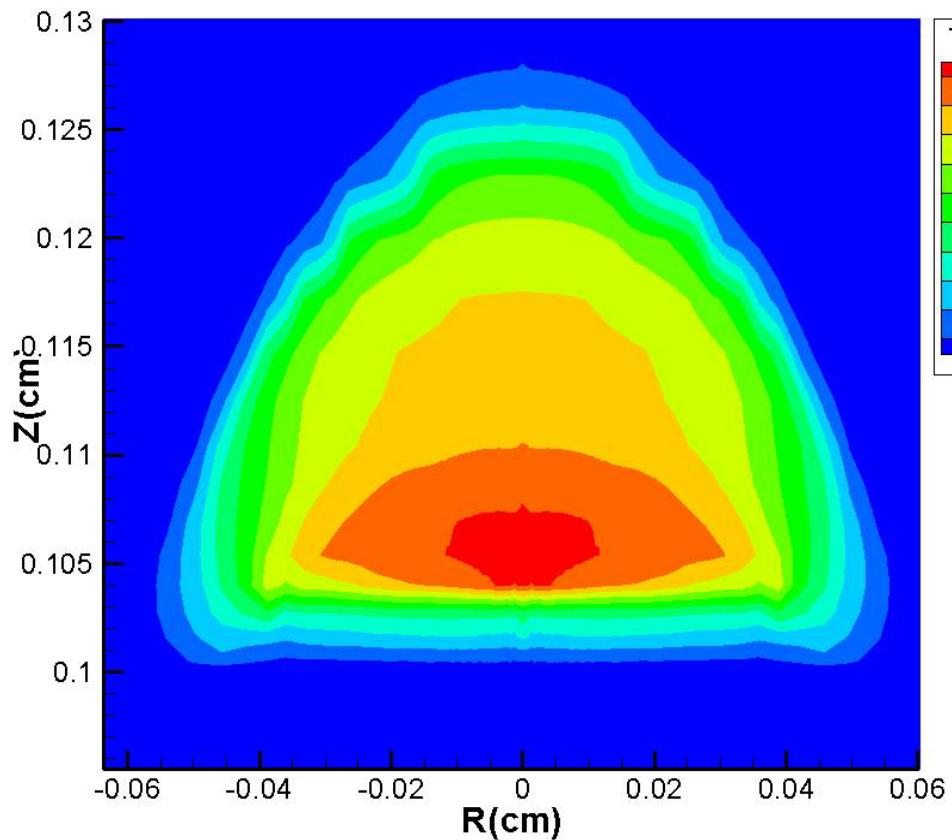


3. 2D plasma: pulse shape (spatial)

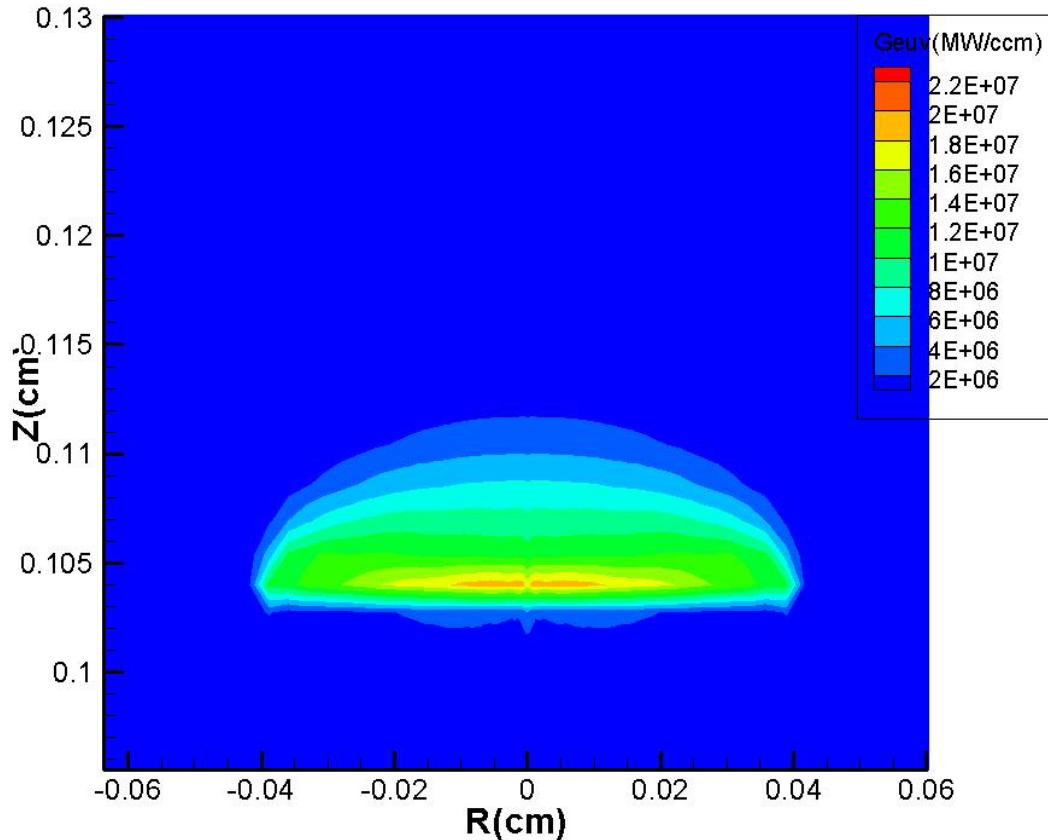
- Z* (EPPRA): 2D LPP simulation
- 2.2-ns **Gaussian** at peak emission



Electron temperature



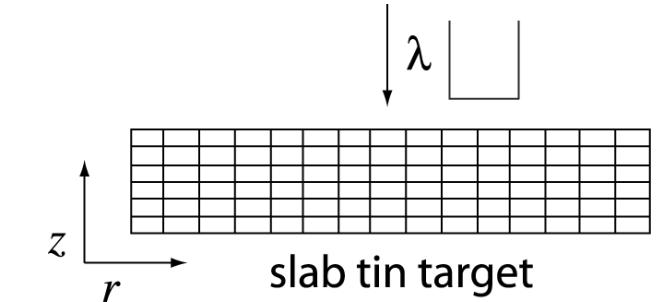
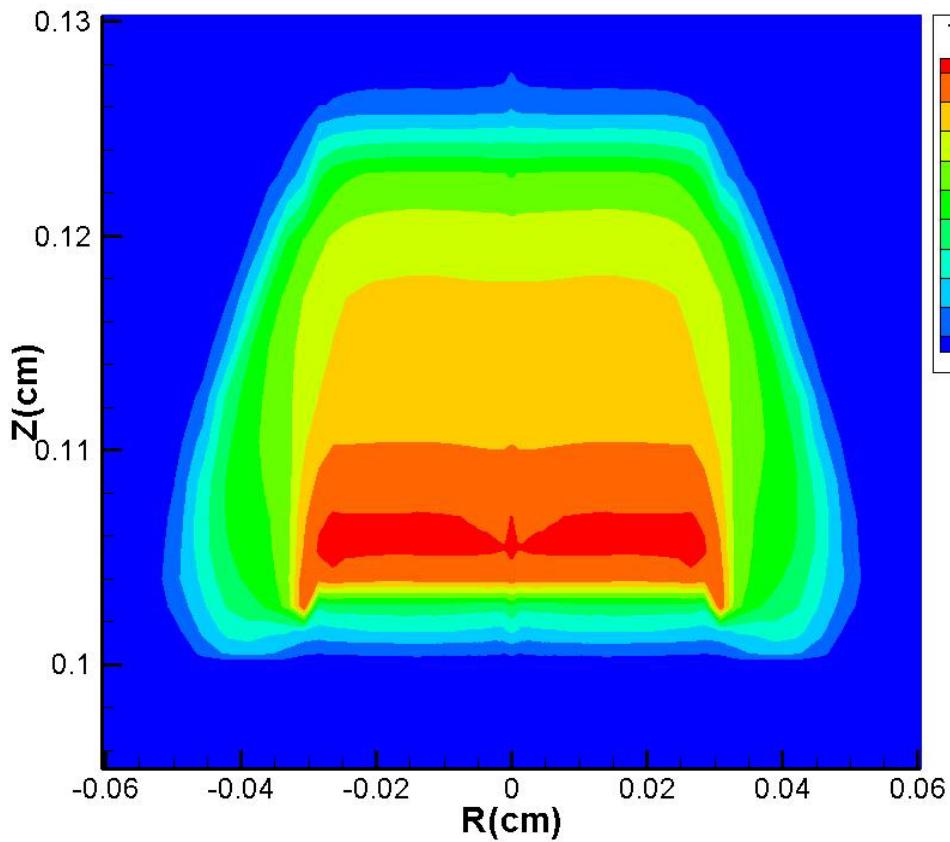
in-band emission (into 4π)



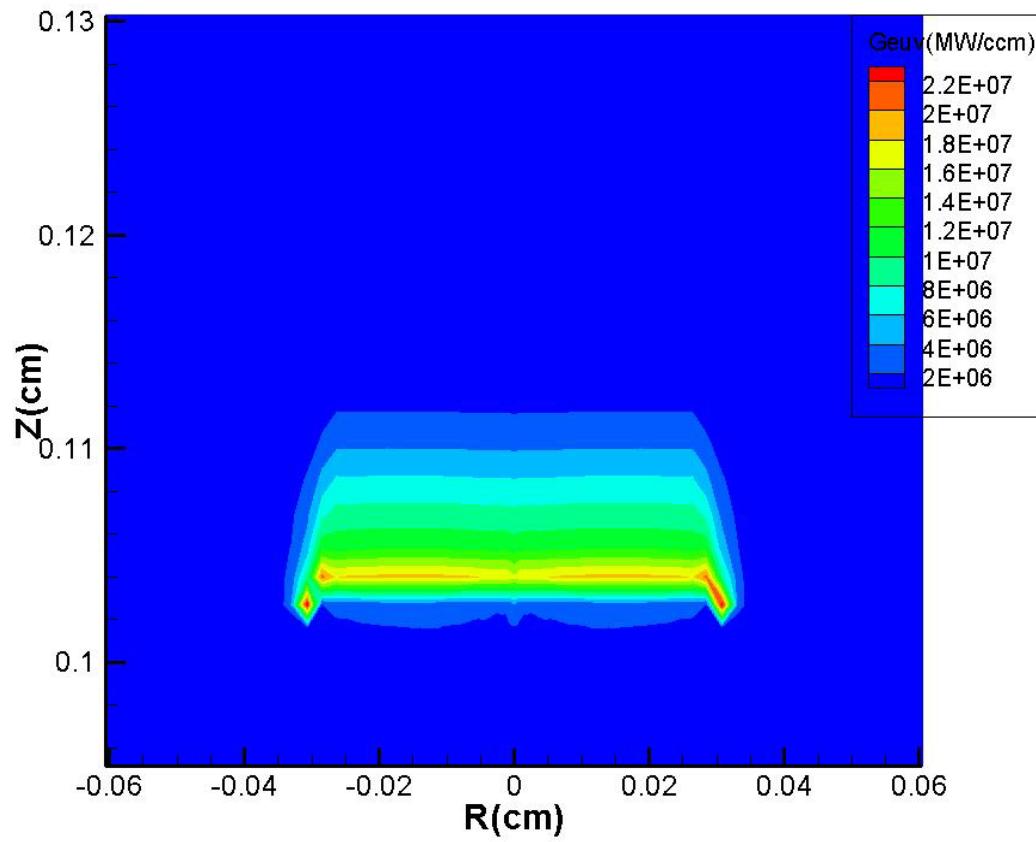
3. 2D plasma: pulse shape (spatial)

- Z* (EPPRA): 2D LPP simulation
- 2.2-ns flat-top at peak emission

Electron temperature



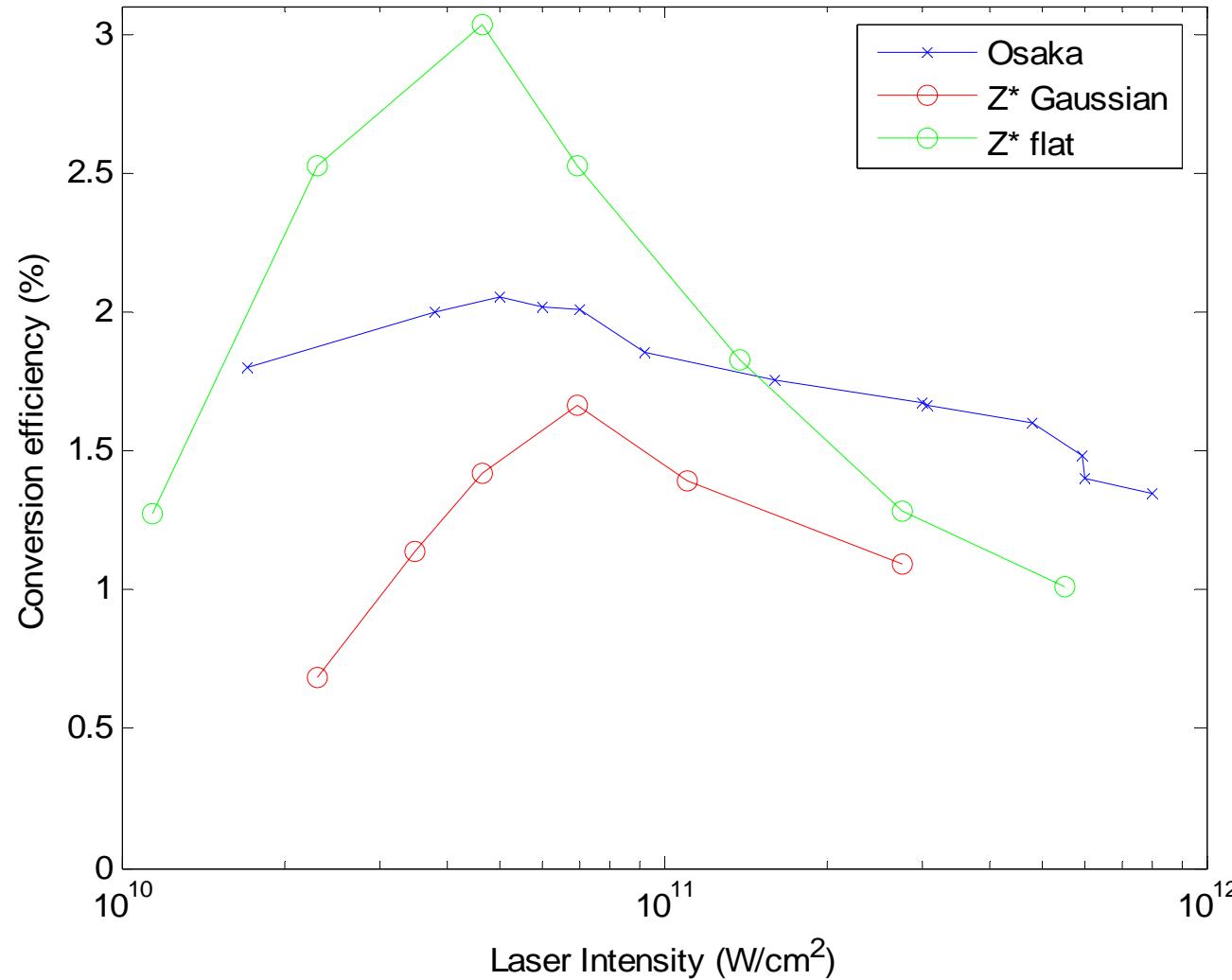
in-band emission (into 4π)



3. 2D plasma: pulse shape (spatial and temporal)

Appl. Phys. Lett., 92, 151501, 2008

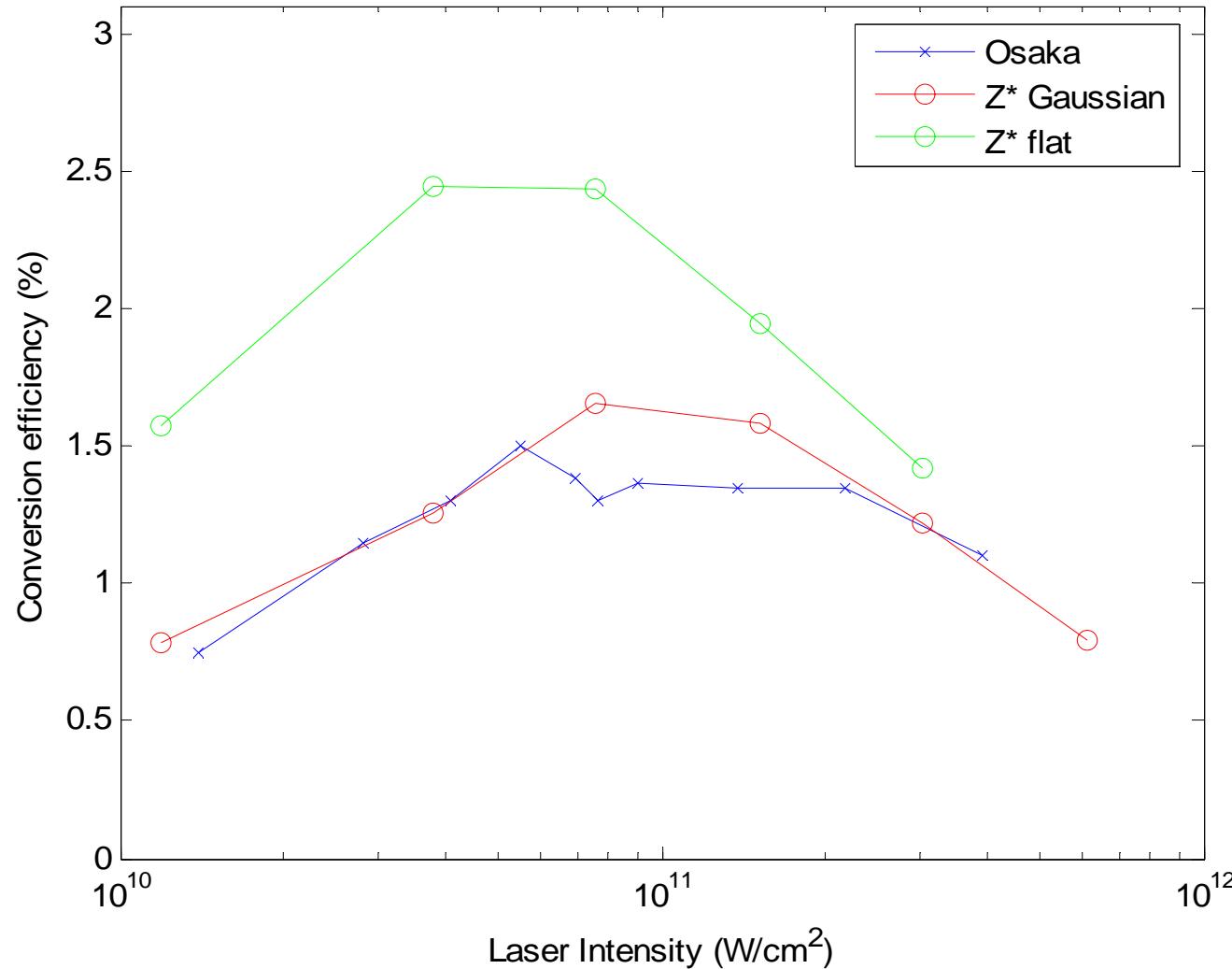
- calculated Gaussian, flat-top, and experiment (Osaka ILE) **2.2 ns**
- CE versus laser power density .66 μm, ..27 μm diameter



3. 2D plasma: pulse shape (spatial and temporal)

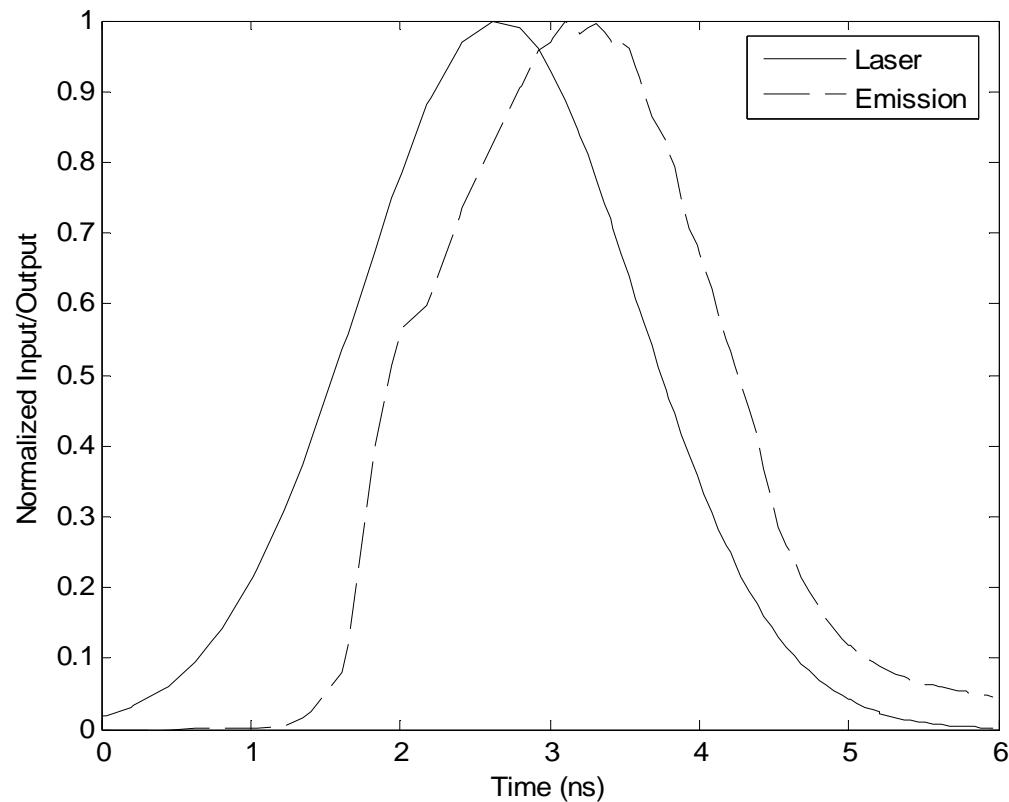
Appl. Phys. Lett., 92, 151501, 2008

- calculated Gaussian, flat-top, and experiment (Osaka ILE) **8.0 ns**
- CE versus laser power density .48 μm, .27 μm diameter

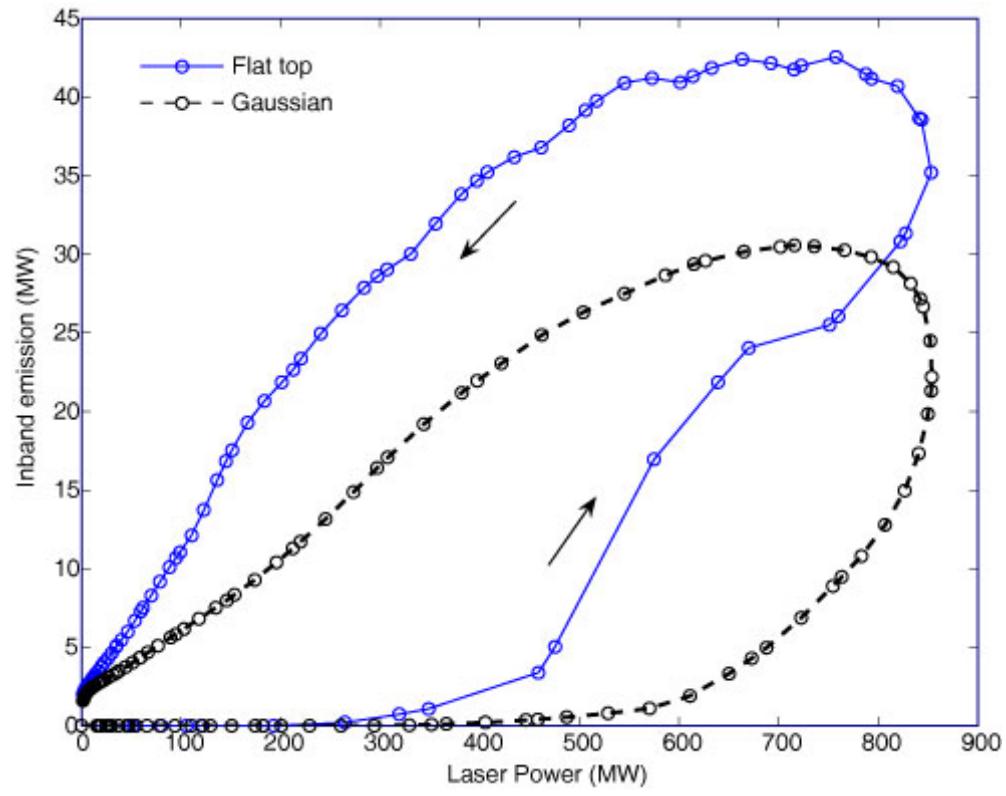


3. 2D plasma (emission lags max pulse)

- Input (laser) and output (spectral response)
- Pulse and emission curves plotted together



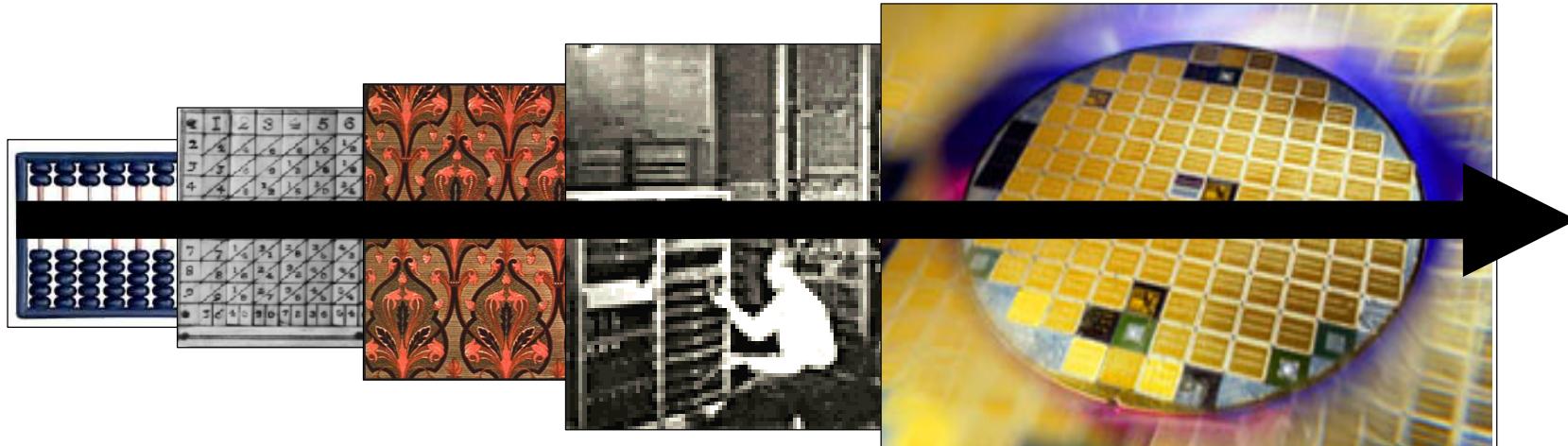
Normalized to one in-band emission and laser power versus time.



In-band emission (into 4π) versus laser power for different beam profiles (flat-top in blue and Gaussian in black).

Conclusions

- steady-state (optically thin) plasma:
 - statistical model: source brightness $F \rightarrow 40$ eV
- time-dependent plasma (optically thick):
 - energy functional for UTAs \rightarrow more manageable computation!
 - radiation transport but no lateral expansion
 - max CE $\rightarrow 0.8 \times 10^{11}$ W/cm² and 10 ns.
- What's next?
 - CO₂ modelling, pulse shaping, and pre-pulse-pulse



Some recent references

- “Tin laser-produced plasma source modelling at 13.5 nm for extreme ultraviolet lithography,” J. White, G. O’Sullivan, S. Zakharov, P. Choi, V. Zakharov, H. Nishimura, S. Fujioka, and K. Nishihara, *Appl. Phys. Lett.*, 92, 151501, 2008
- “Optimising 13.5-nm laser-produced tin plasma emission as a function of laser wavelength,” J. White, P. Dunne, P. Hayden, F. O’Reilly, and G. O’Sullivan, *Appl. Phys. Lett.*, 90, 181502, 2007.
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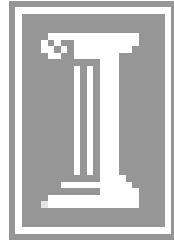
UCD Modelling Collaborations



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突破EUV光源的瓶颈

Chenzong Dong, Mao Gen Su



David Ruzic



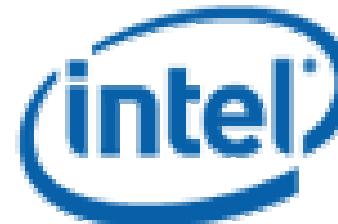
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